

Modelling of bulk superconductor magnetisation

Dr Mark Ainslie

Royal Academy of Engineering (UK) Research Fellow

**Co-authors: Prof. Hiroyuki Fujishiro, Iwate University
Jin Zou, University of Cambridge**

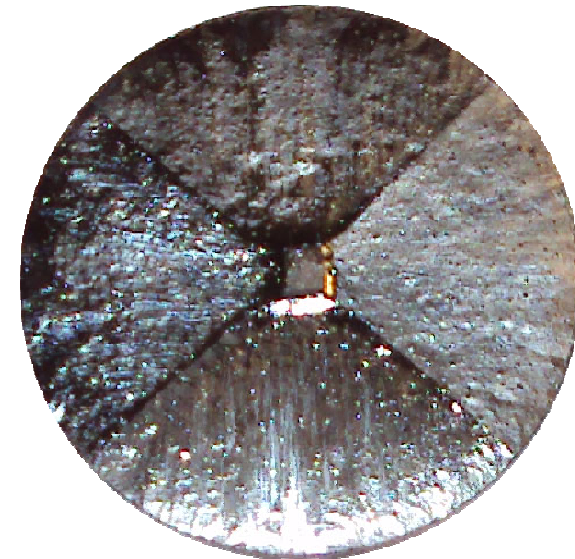
Bulk Superconductivity Group, Department of Engineering

Presentation Outline

- **Overview of numerical modelling of bulk superconductor magnetisation**
- **Case study #1: Field cooling magnetisation of bulk MgB_2**
- **Case study #2: Split coil, pulsed field magnetisation of bulk high-temperature superconductors with an iron yoke**

Bulk Superconductors

- Bulk superconducting materials trap magnetic flux via macroscopic electrical currents
- Magnetisation increases with sample volume



A large, single grain bulk superconductor

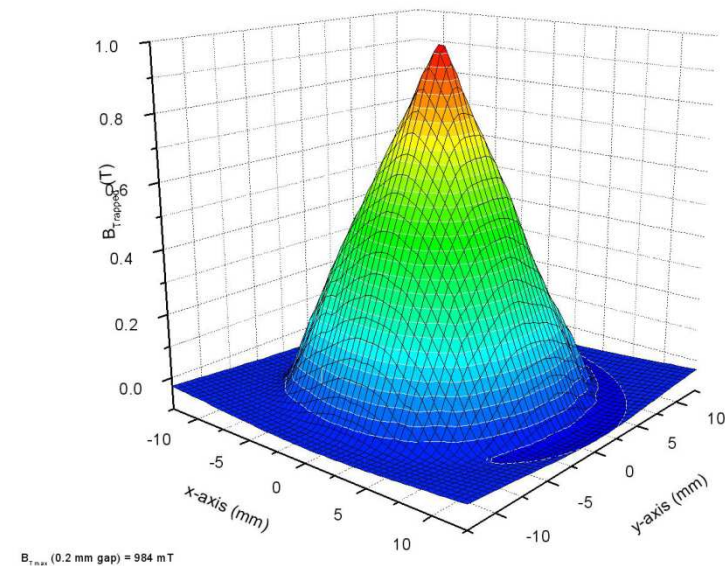
Bulk Superconductors

- Bulk superconducting materials trap magnetic flux via macroscopic electrical currents
 - Magnetisation increases with sample volume
- Trapped field given by

$$B_{\text{trap}} = k \mu_0 J_c R$$

$$k = \frac{t_B}{2R} \cdot \ln \left(\frac{R + \sqrt{R^2 + t_B^2}}{t_B} \right)$$

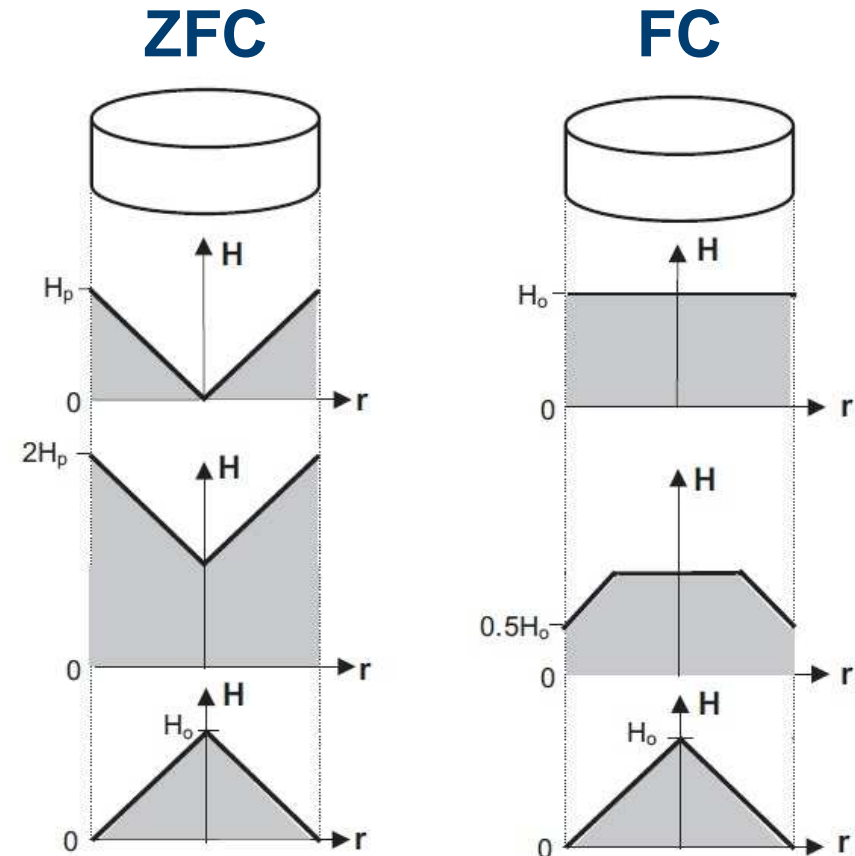
from Biot-Savart law + Bean model (infinite slab) = constant J_c



Typical trapped magnetic field profile of a bulk superconductor

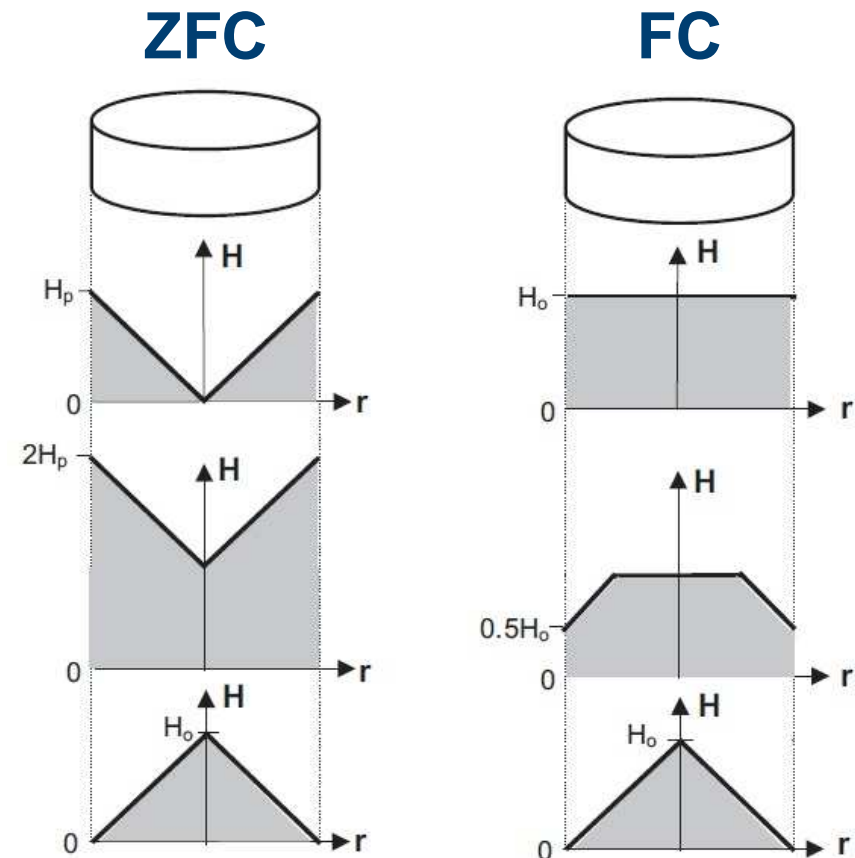
Magnetisation of Bulk Superconductors

- Three magnetisation techniques:
 - Field Cooling (FC)
 - Zero Field Cooling (ZFC)
 - Pulsed Field Magnetisation (PFM)



Magnetisation of Bulk Superconductors

- Three magnetisation techniques:
 - Field Cooling (FC)
 - Zero Field Cooling (ZFC)
 - Pulsed Field Magnetisation (PFM)
- To trap B_{trap} , need at least B_{trap} or higher
 - FC and ZFC require large magnetising coils
 - Impractical for applications/devices



Pulsed Field Magnetisation

- **PFM technique: compact, mobile, relatively inexpensive**
- **Issues: B_{trap} [PFM] < B_{trap} [FC], [ZFC]**
 - Temperature rise ΔT due to rapid movement of magnetic flux
- **Record PFM trapped field: 5.2 T @ 29 K** (45 mm diameter Gd-Ba-Cu-O)
Fujishiro et al. *Physica C* 2006
 - Record trapped field by FC: 17.6 T @ 26 K (2 x 25 mm dia Gd-Ba-Cu-O)
Durrell et al. *Supercond. Sci. Technol.* 2014
- **Many considerations:**
 - Pulse magnitude, pulse duration, temperature, number of pulses, shape of magnetising coil(s)
 - Dynamics of magnetic flux during PFM process

Numerical Modelling of Magnetisation

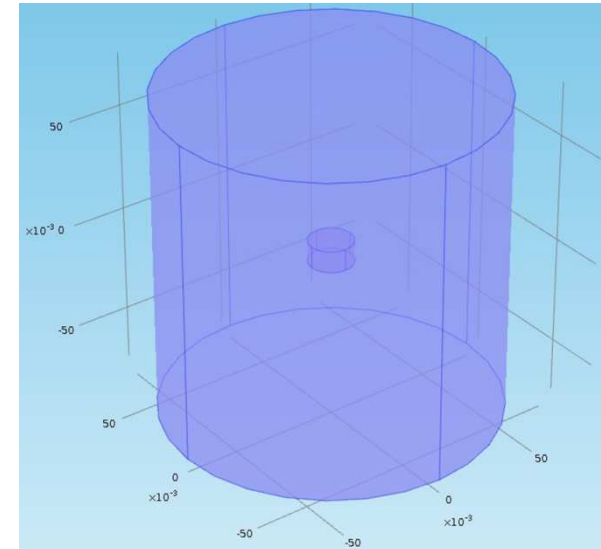
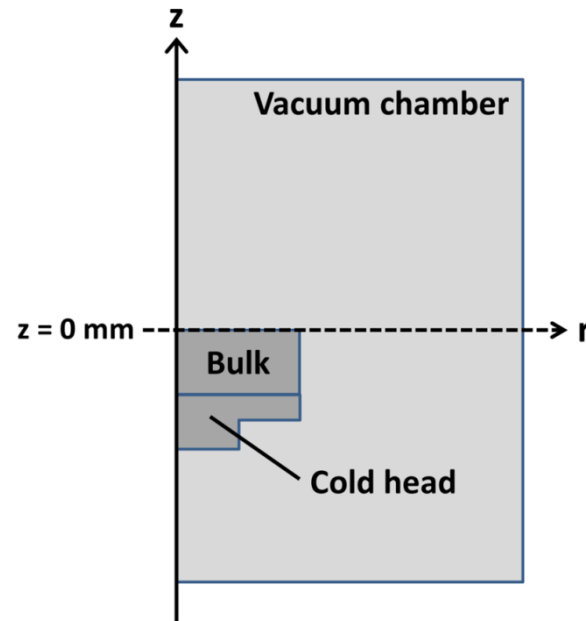
BULK GEOMETRY &
MAGNETISATION
FIXTURE

ELECTROMAGNETIC
FORMULATION

THERMAL
EQUATIONS &
PROPERTIES

$J_c(B, T)$

E-J POWER LAW



- 2D axisymmetric generally sufficient for cylindrical bulks with a homogeneous J_c distribution
- 3D required for an inhomogeneous J_c distribution around the ab -plane; for non-symmetric shapes

Numerical Modelling of Magnetisation

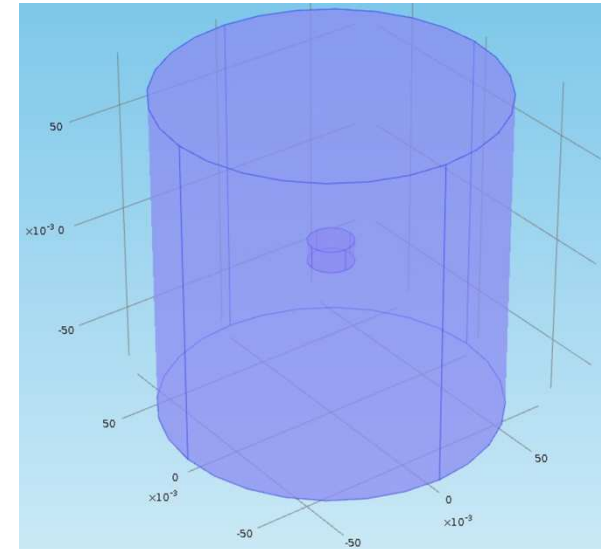
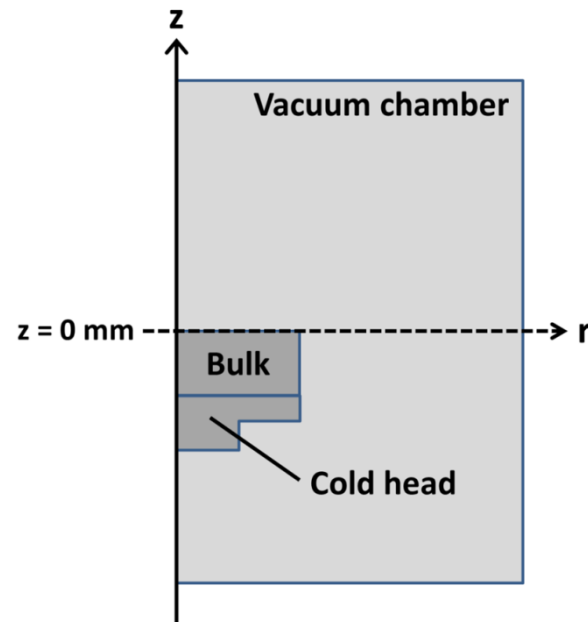
BULK GEOMETRY &
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- Magnetising fixture can be simulated by using uniform boundary conditions or by inserting a copper coil sub-domain
- Cooling can be simulated using a cold head / vacuum chamber (left) or by submersion in liquid cryogen (right)

Numerical Modelling of Magnetisation

BULK GEOMETRY &
MAGNETISATION
FIXTURE

Finite element method is commonly used & well developed (other techniques do exist)

ELECTROMAGNETIC
FORMULATION

Governing equations:

Maxwell's equations (*H* formulation)

THERMAL
EQUATIONS &
PROPERTIES

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} = -\frac{\partial(\mu_0 \mu_r \mathbf{H})}{\partial t}$$

Faraday's Law

$$\nabla \times \mathbf{H} = \mathbf{J}$$

Ampere's Law

$J_c(B, T)$

Other formulations also exist (*A-V*, *T-Ω*, *Campbell's equation*)

E-J POWER LAW

Numerical Modelling of Magnetisation

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Thermal behaviour needs to be modelled when the bulk experiences a significant change in temperature, e.g., during PFM or modelling complete FC magnetisation process

Governing equations:

$$\rho \cdot C \frac{dT}{dt} = \nabla \cdot (k \nabla T) + Q$$

$$Q = E_{\text{norm}} \cdot J_{\text{norm}}$$

ρ = mass density, C = specific heat, κ = thermal conductivity, Q = heat source

Numerical Modelling of Magnetisation

BULK GEOMETRY &
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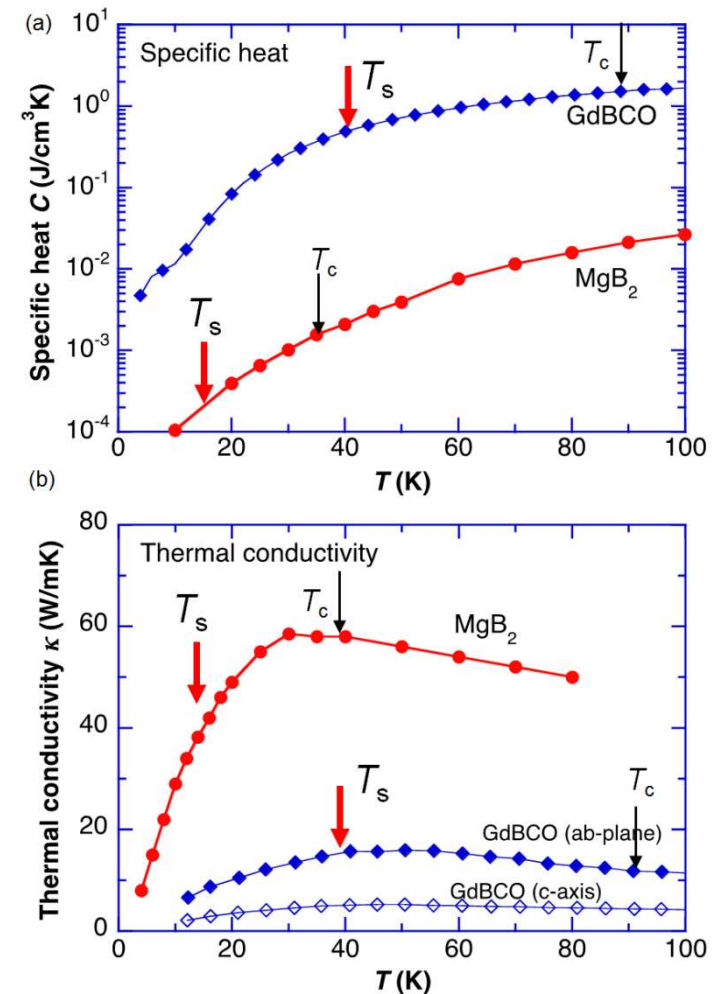
THERMAL
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$J_c(B, T)$

E-J POWER LAW

Can choose constant parameters for C , κ for $T = T_{op}$ as a reasonable approximation

Can use measured experimental data/fitting function over a specific temperature range



Numerical Modelling of Magnetisation

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(RE)BCO materials

Kim-like model:

$$J_c = \frac{J_{c0}}{\left(1 + \frac{B}{B_0}\right)^\alpha}$$

Fishtail effect: $J_c(B) = J_{c1} \exp\left(-\frac{B}{B_L}\right) + J_{c2} \frac{B}{B_{\max}} \exp\left[\frac{1}{y} \left(1 - \left(\frac{B}{B_{\max}}\right)^y\right)\right]$

MgB₂ materials

$$J_c(B, T) = J_{c0}(T) \exp\left(-\frac{B}{B_0}\right)^a$$

- Some software packages, e.g., COMSOL, allow direct interpolation of experimental data without need for data fitting

e.g., Hu et al. *Supercond. Sci. Technol.* **28** (2015) 065011

- Important for coated conductor modelling where in-field behaviour can be quite complex: $J_c(B, \theta, T)$

Numerical Modelling of Magnetisation

BULK GEOMETRY &
MAGNETISATION
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THERMAL
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PROPERTIES

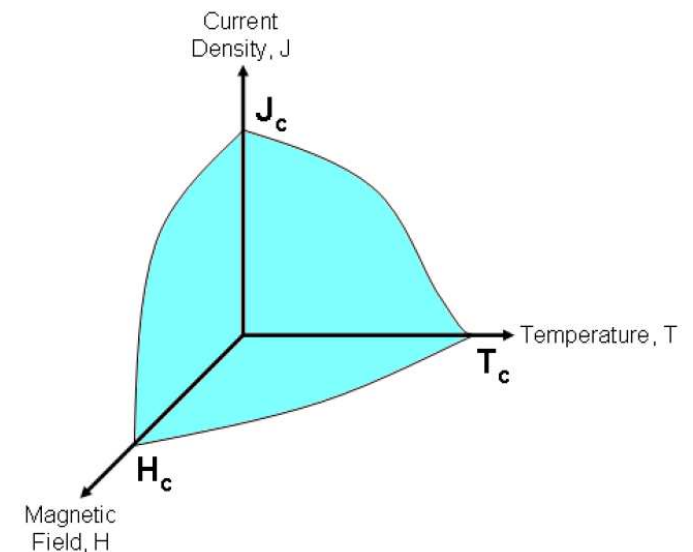
$J_c(B, T)$

E-J POWER LAW

Linear $J_{c0}(T)$ relationship
has also been used in the
literature

Assumption made is that
in-field behaviour,
 $J_c(B, T=T_{op})$, doesn't
change for variations
around the operating
temperature, T_{op}

$$J_{c0}(T) = \alpha \left[1 - \left(\frac{T}{T_c} \right)^2 \right]^{1.5}$$



Numerical Modelling of Magnetisation

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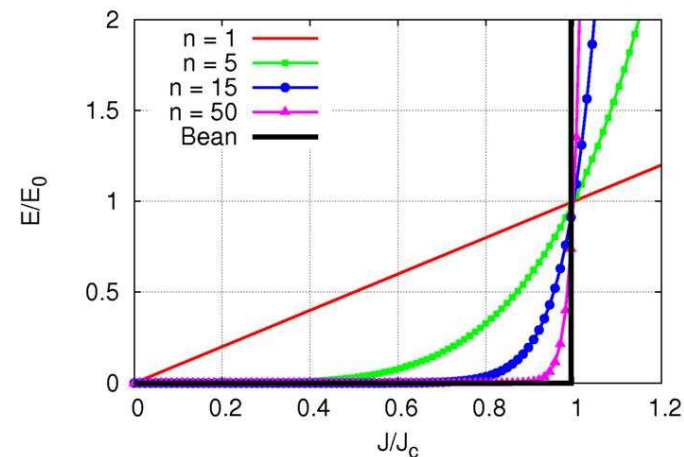
THERMAL
EQUATIONS &
PROPERTIES

$J_c(B, T)$

E-J POWER LAW

***E*-*J* power law**

- Conventional conductors → non-linear permeability, linear resistivity
- Superconductors → linear permeability (μ_0), non-linear resistivity
- Non-linearity is extreme: power law with $n > 20$



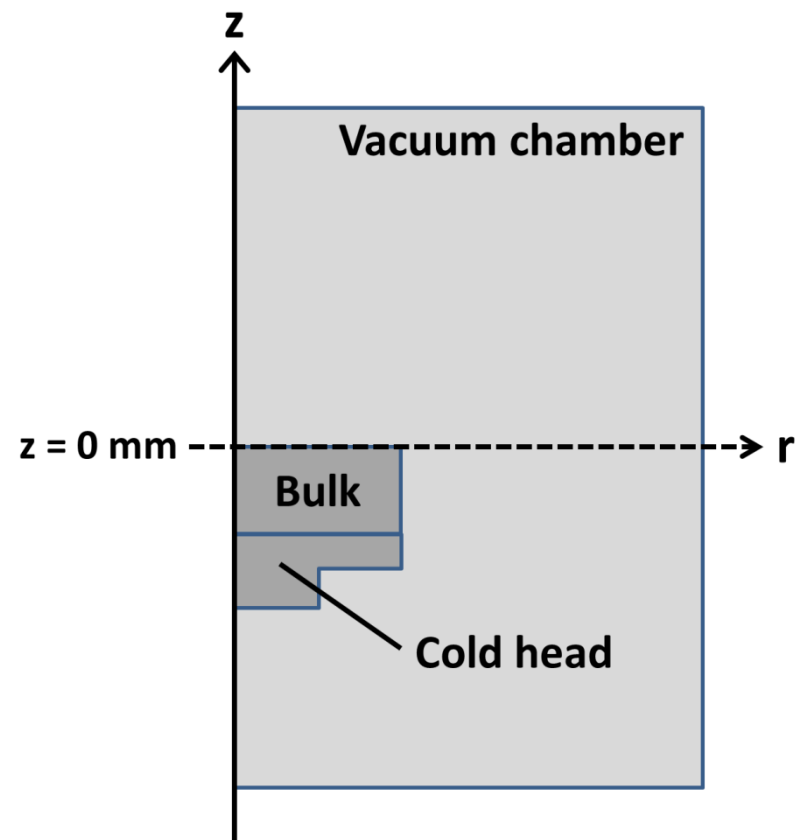
$$\mathbf{E} = E_0 \left(\frac{J}{J_c} \right)^{n-1} \frac{J}{J_c}$$

Case Study:

Field Cooling Magnetisation of MgB_2 Bulk Superconductors

Bulk MgB₂ Modelling

- Numerical modelling of bulk MgB₂ is simpler than (RE)BCO
 - J_c distribution generally quite homogeneous
 - 2D axisymmetric geometry, so $\mathbf{H} = [H_r, H_z]$, $\mathbf{J} = [J_\phi]$, $\mathbf{E} = [E_\phi]$
 - Can use measured $J_c(B, T)$ characteristics of a single, small specimen



Bulk MgB₂ Modelling – $J_c(B, T)$ Data Fitting

- **Four samples measured:**
 - Trapped field (FC) between ~ 5-15 K and 40 K
 - $J_c(B, T)$ of single, small specimen

Bulk MgB₂ sample information.

Sample	HIP#22	HIP#38	HIP-Ti20%	IG1
T_c (K)	38.5	38.5	39	37.5
Diameter, d (mm)	22	38	36	32
Thickness, t_B (mm)	18	7	7	6
Aspect ratio (d/t_B)	1.2	5.4	5.1	5.33
Relative mass density	93%	93%	94%	90%
Maximum trapped field, B_z (FC at 20 K) (T)	1.92	2.09	2.77	1.34
J_c (FC at 20 K) (A m ⁻²) ^a	2.94×10^8	2.76×10^8	4.23×10^8	2.08×10^8
Reference	[11]	[11]	[11]	[22]

Zou, Ainslie, Fujishiro et al. *Supercond. Sci. Technol.* **28** (2015) 075009

Bulk MgB₂ Modelling – Field Cooling Magnetisation

- **Simulating FC magnetisation process:**

- FC with $B_{\text{app}} = B_{\text{trap}}$:

1. $0 \leq t \leq x_1$ Apply ramped field to $B_{\text{app}} = B_{\text{trap}}$ at $T = T_{\text{ex}} > T_c$
2. $x_1 \leq t \leq x_2$ Slow cooling of bulk to operating temperature, $T = T_{\text{op}}$
3. $x_2 \leq t \leq x_3$ Slowly ramp applied field $B_{\text{ex}} \rightarrow 0$

- In electromagnetic model, need to define ρ for all temperatures:

- For $T > T_c$, need to define ρ_{normal} ($\rho_{\text{normal}} = 3 \times 10^{-8} \Omega\text{m}$)

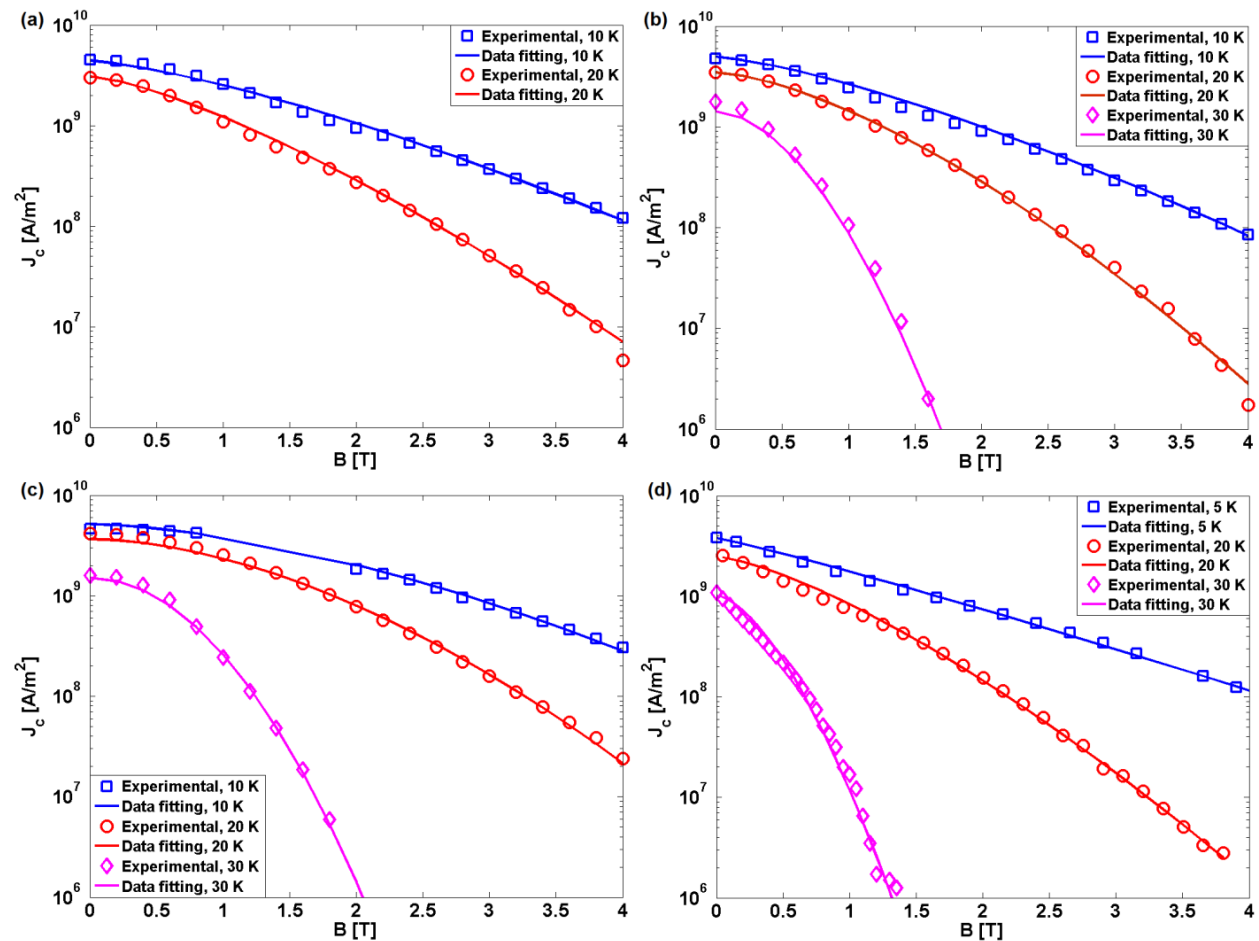
- For $T < T_c$, ρ_{sc} defined from E - J power law, where $\mathbf{E} = \rho \mathbf{J}$: $\mathbf{E} = E_0 \left(\frac{J}{J_c} \right)^{n-1} \frac{\mathbf{J}}{J_c}$

- To avoid non-convergence at T_c : $\rho(T) = \frac{\rho_{\text{sc}} \cdot \rho_{\text{normal}}}{\rho_{\text{sc}} + \rho_{\text{normal}}}$

Zou, Ainslie, Fujishiro et al. *Supercond. Sci. Technol.* **28** (2015) 075009

Bulk MgB₂ Modelling – J_c(B, T) Data Fitting

$$J_c(B, T) = J_{c0}(T) \exp\left(-\frac{B}{B_0}\right)^a$$



Zou, Ainslie, Fujishiro et al. *Supercond. Sci. Technol.* **28** (2015) 075009

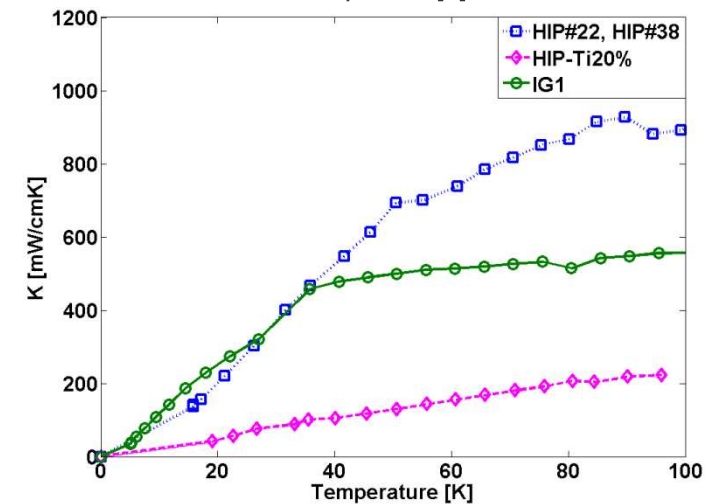
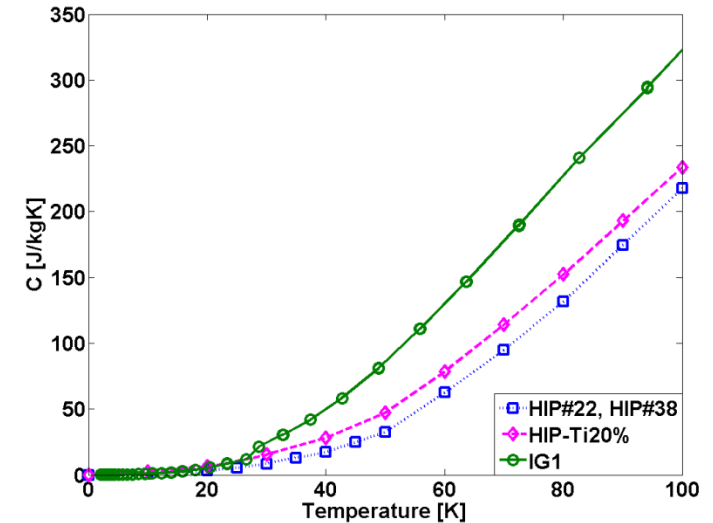
Bulk MgB₂ Modelling – Thermal Properties

Thermal equation: $\rho \cdot C \frac{dT}{dt} = \nabla \cdot (k \nabla T) + Q$

- Here T changes from $T_{\text{ex}} = 100 \text{ K}$ ($> T_c$) to $T_{\text{op}} = 5 - 30 \text{ K}$
- Measured experimental data from 0 – 100 K for each sample input directly into model (direct interpolation)

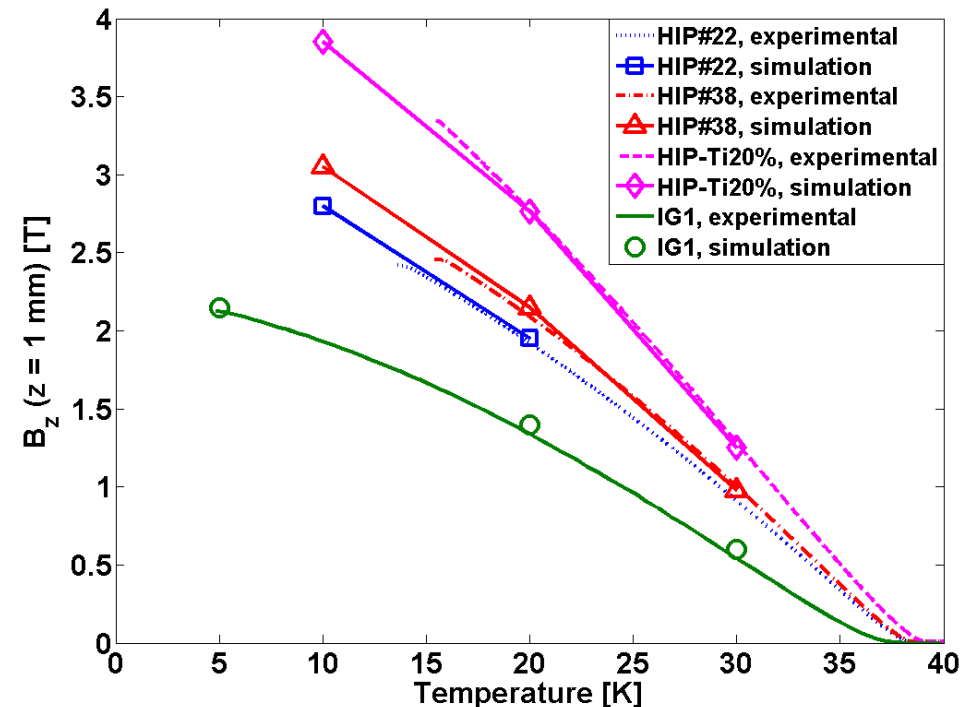
$$J_c(B, T) = J_{c0}(T) \exp\left(-\frac{B}{B_0}\right)^a$$

$$J_{c0}(T) = \alpha \left[1 - \left(\frac{T}{T_c} \right)^2 \right]^{1.5}$$



Bulk MgB₂ Modelling – Comparison of Results

- **Simulation reproduces experimental trapped field measurements extremely well**
- Samples have excellent homogeneity
- Model is validated as a fast & accurate tool to predict trapped field performance
 - Any size of bulk MgB₂ disc & any specific operating conditions
 - Can investigate current distribution / trapped field in partially magnetised situation



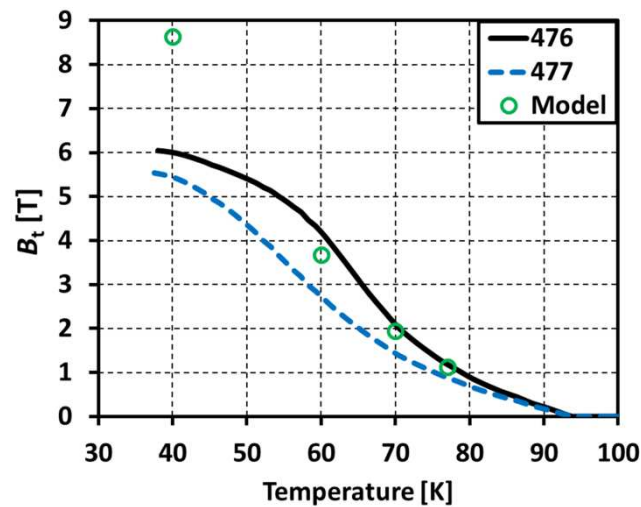
Comparison of simulation & experimental results for trapped field in different bulk MgB₂ samples

Zou, Ainslie, Fujishiro et al. *Supercond. Sci. Technol.* **28** (2015) 075009

Case Study #2:

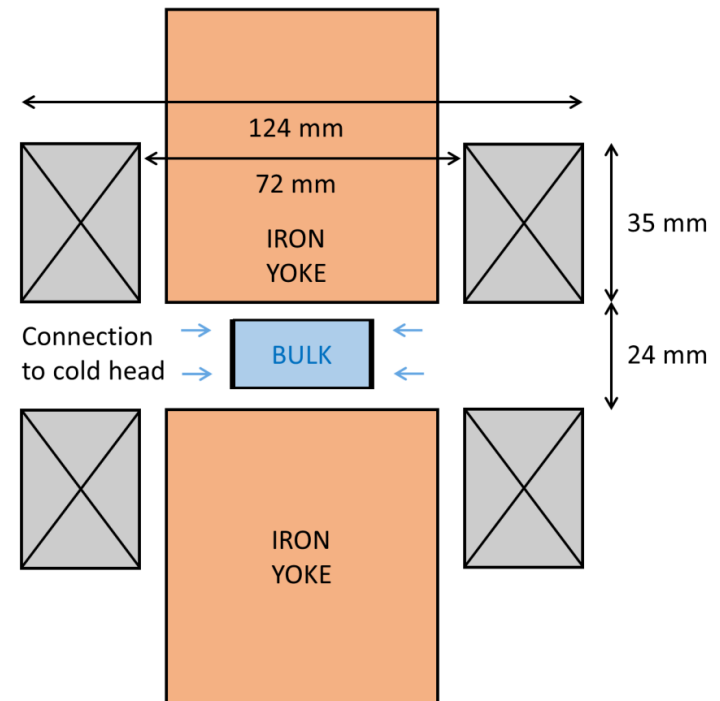
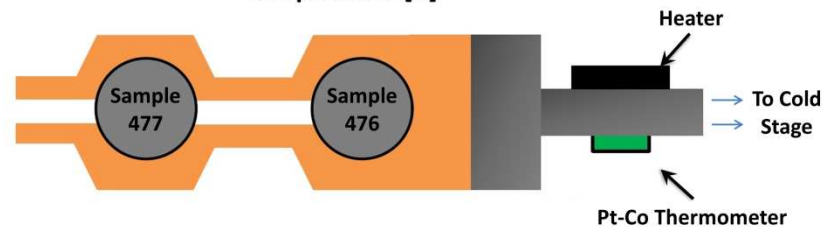
Split Coil, Pulsed Field Magnetisation of
Bulk High-Temperature Superconductors
With an Iron Yoke

Split Coil PFM with an Iron Yoke



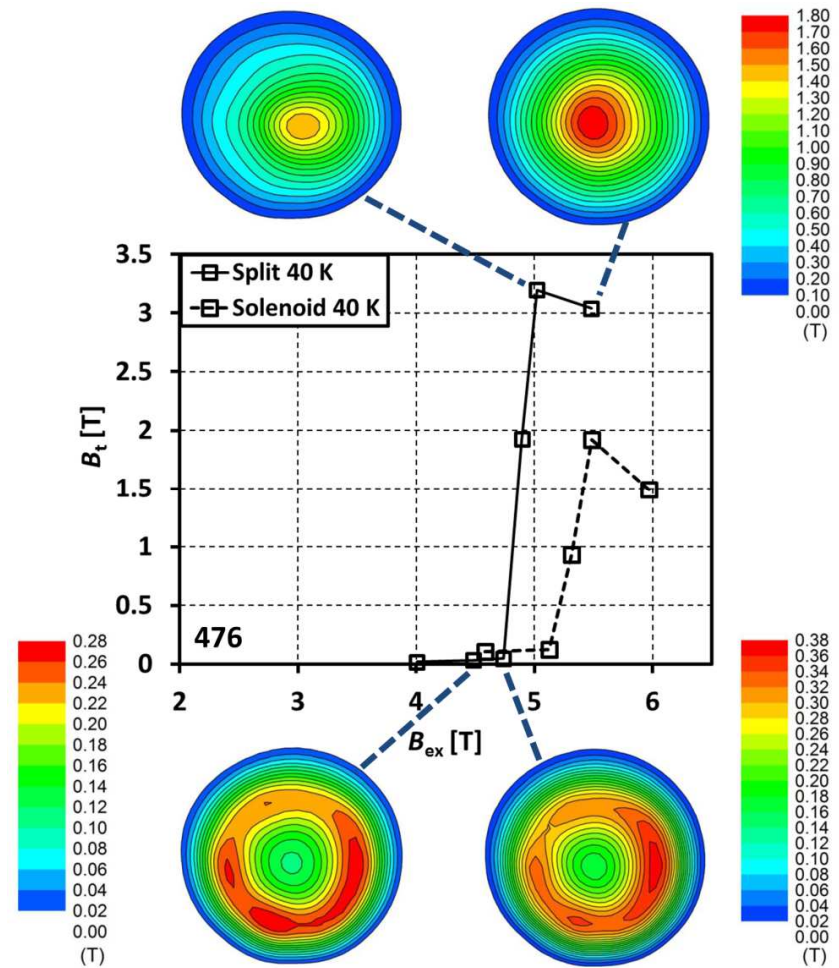
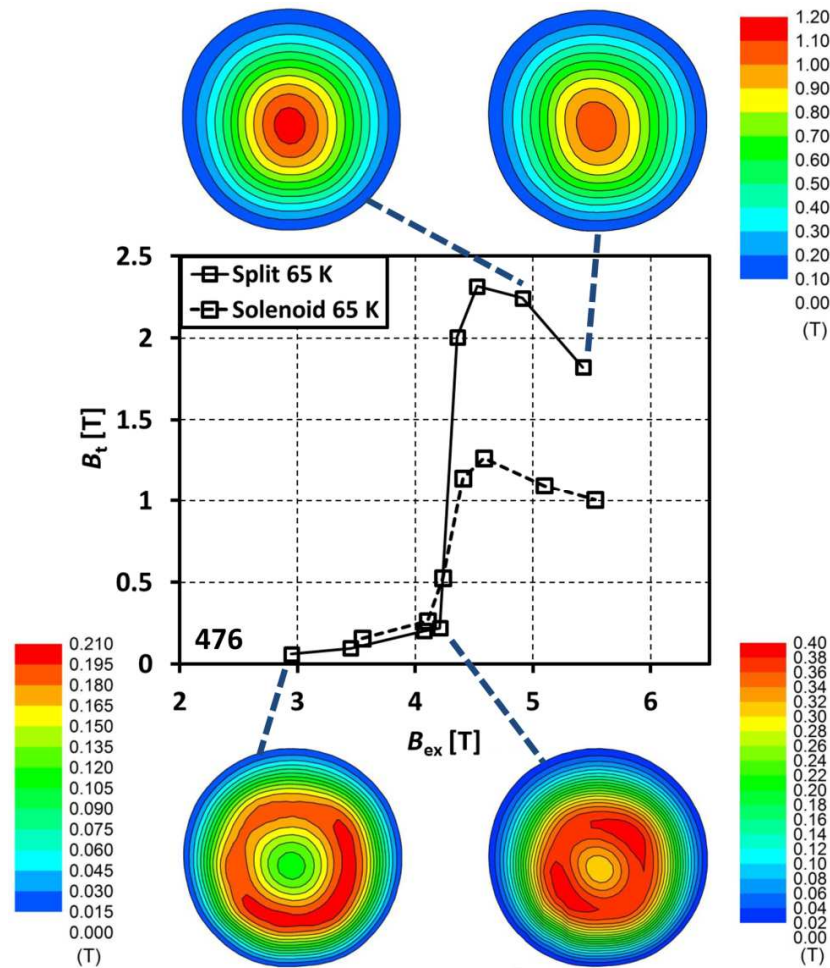
Two Ag-containing (15 wt%) Gd-Ba-Cu-O samples:
 30 mm diameter, 15 mm thickness

#476 sample $B_t = 6$ T (40 K) 3.11 T (65 K)
 #477 sample $B_t = 5.44$ T (40 K) 2.02 T (65 K)



Ainslie, Fujishiro et al. *Supercond. Sci. Technol.* **29** (2016) 074003

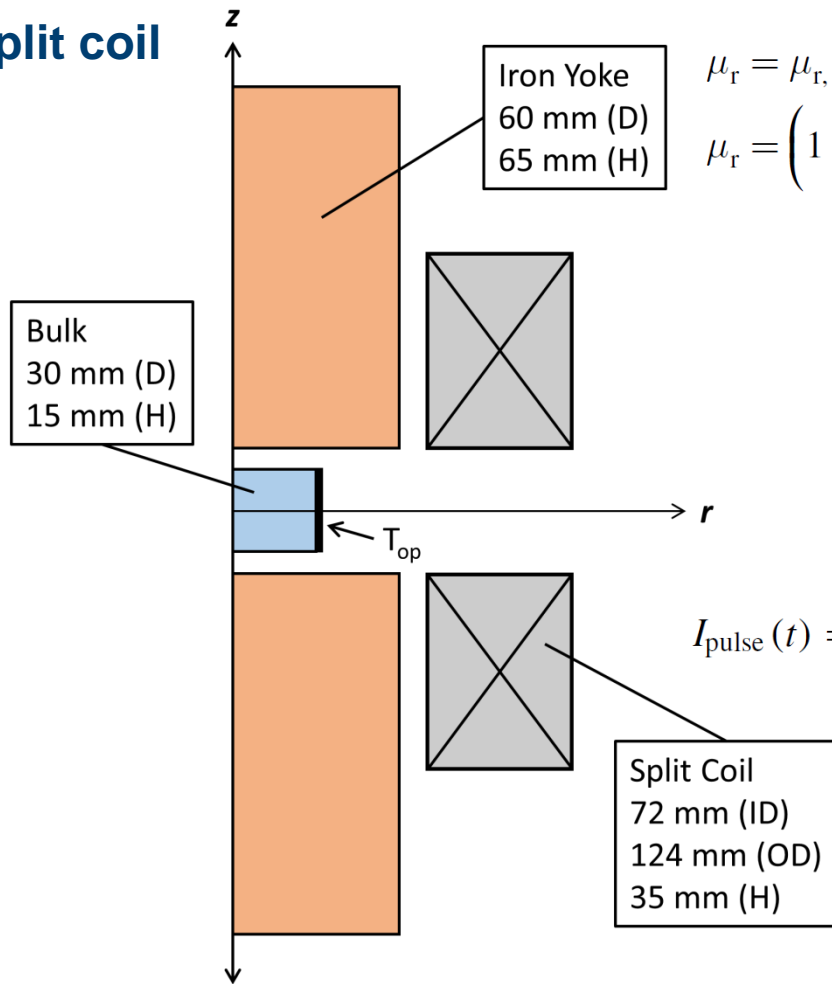
Split Coil PFM with an Iron Yoke



Ainslie, Fujishiro et al. *Supercond. Sci. Technol.* **29** (2016) 074003

Split Coil PFM with an Iron Yoke

Split coil

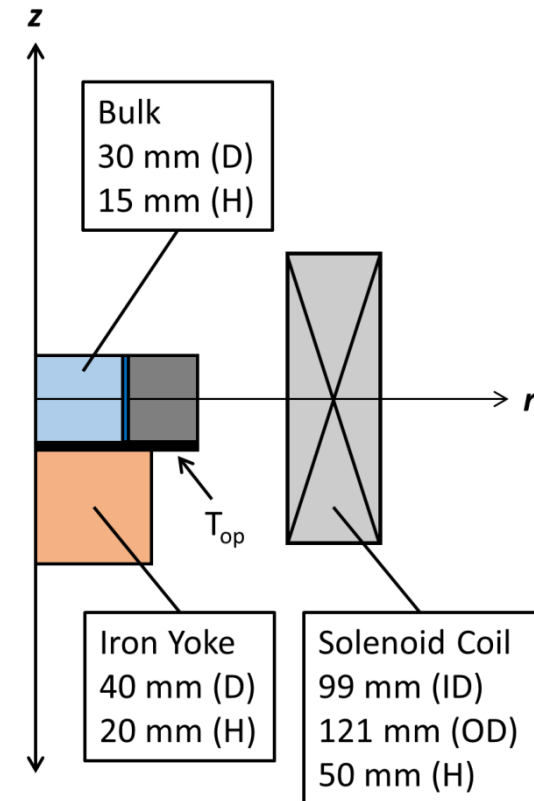


$$\mu_r = \mu_{r, \max} \text{ for } B < \mu_0 M_{\text{sat}}$$

$$\mu_r = \left(1 + \frac{M_{\text{sat}}}{H} \right) \text{ for } B \geq \mu_0 M_{\text{sat}}$$

$$I_{\text{pulse}}(t) = N \cdot I_0 \frac{t}{\tau} \exp\left(1 - \frac{t}{\tau}\right)$$

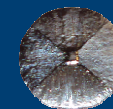
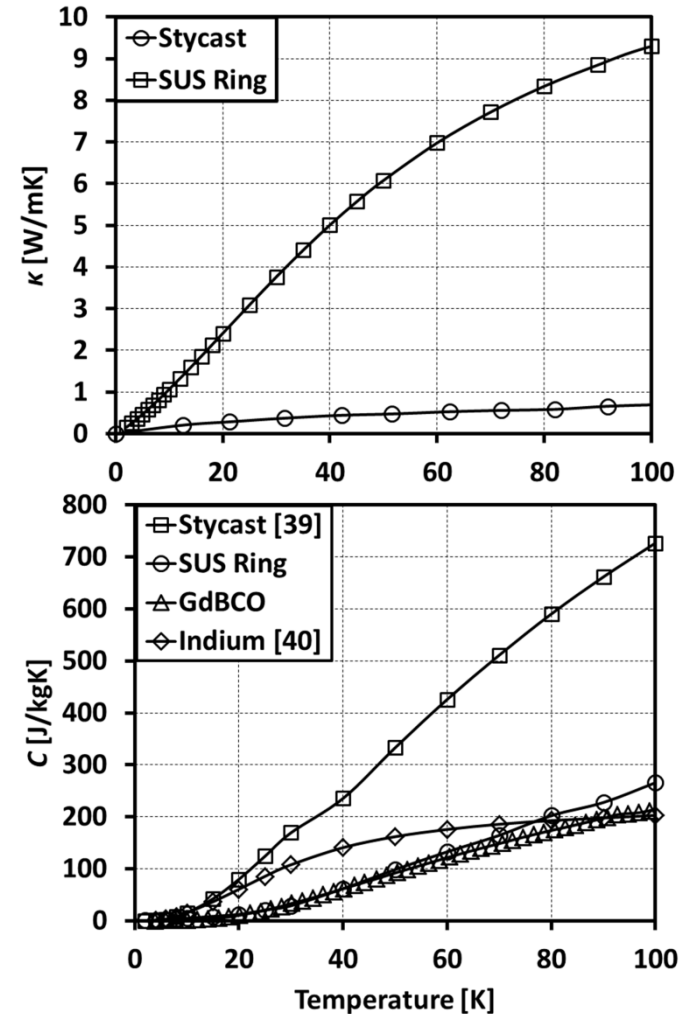
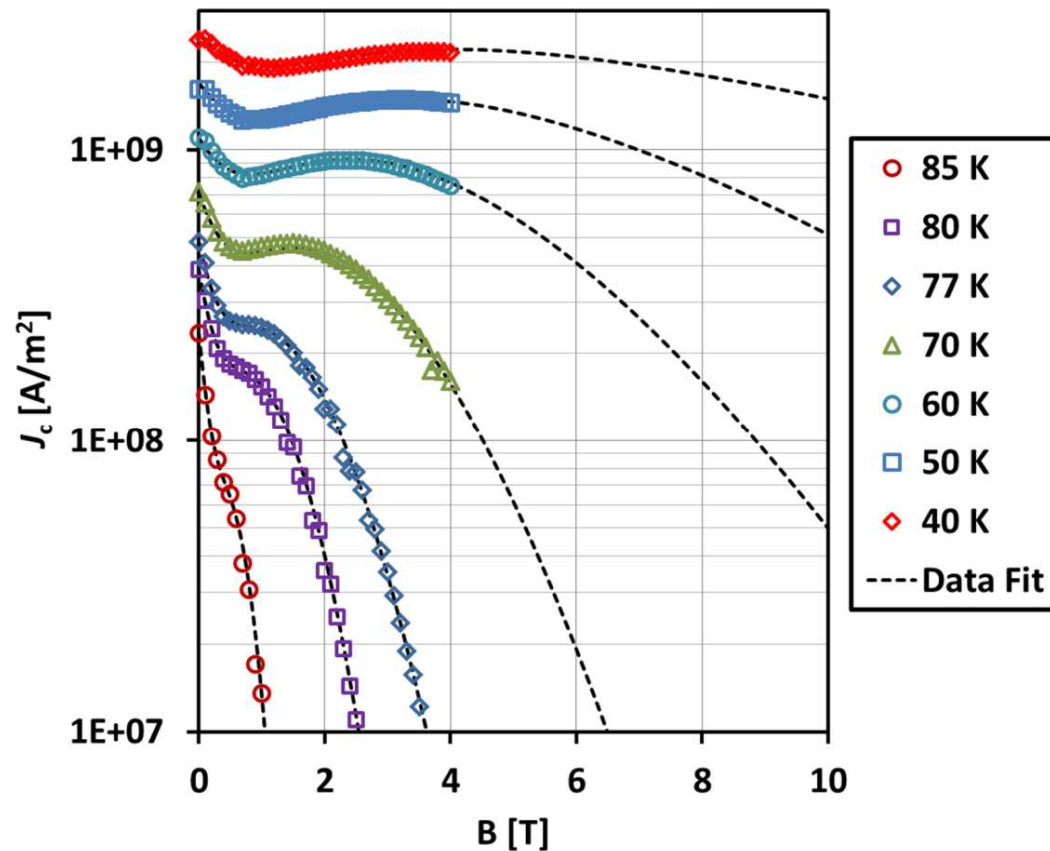
Solenoid coil



Ainslie, Fujishiro et al. *Supercond. Sci. Technol.* **29** (2016) 074003

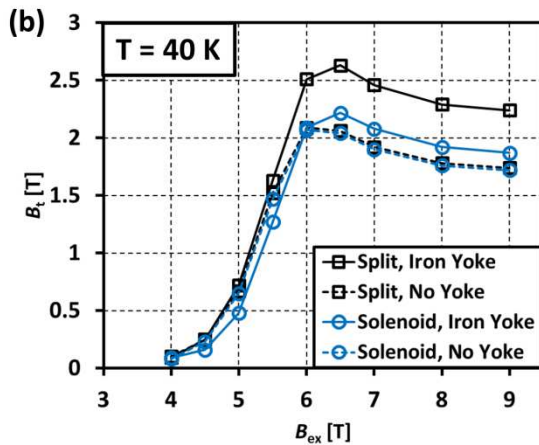
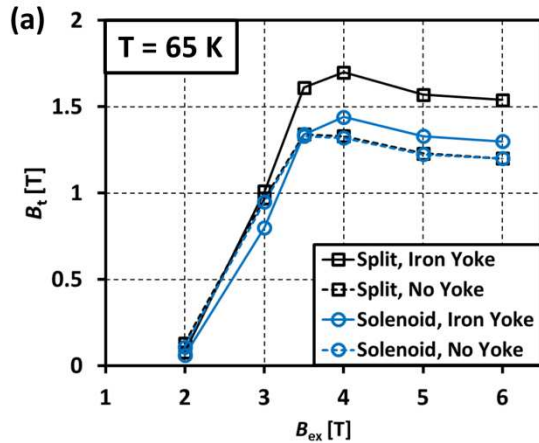
Split Coil PFM with an Iron Yoke

$$J_c(B) = J_{c1} \exp\left(-\frac{B}{B_L}\right) + J_{c2} \frac{B}{B_{\max}} \exp\left[\frac{1}{y} \left(1 - \left(\frac{B}{B_{\max}}\right)^y\right)\right]$$

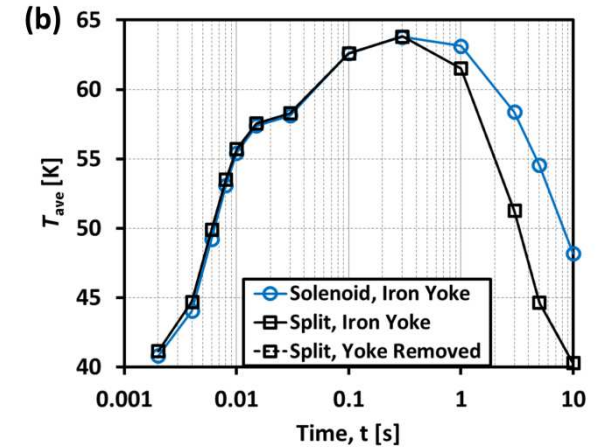
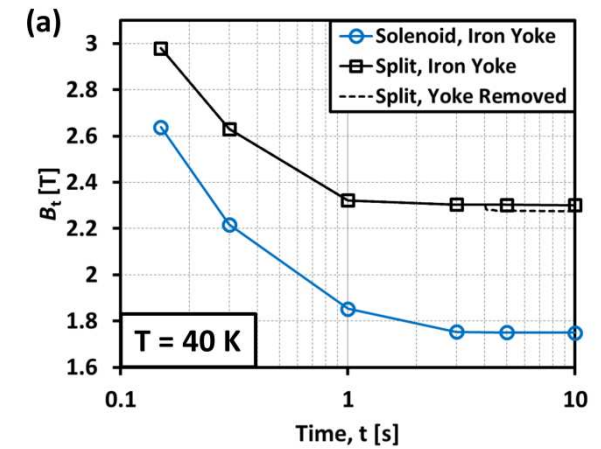
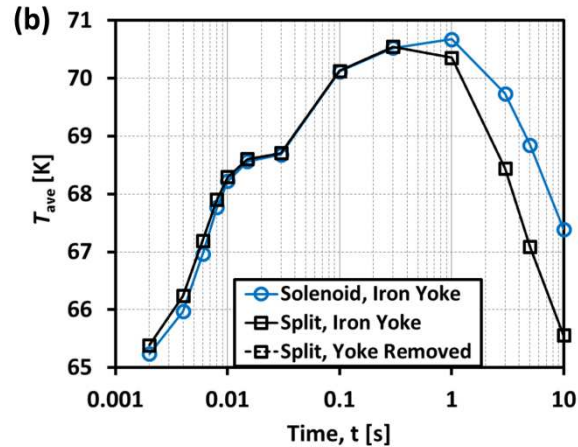
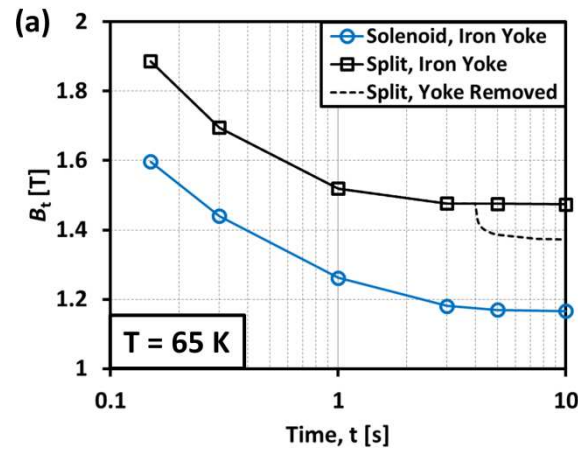


Split Coil PFM with an Iron Yoke

Trapped field @ $t = 300$ ms



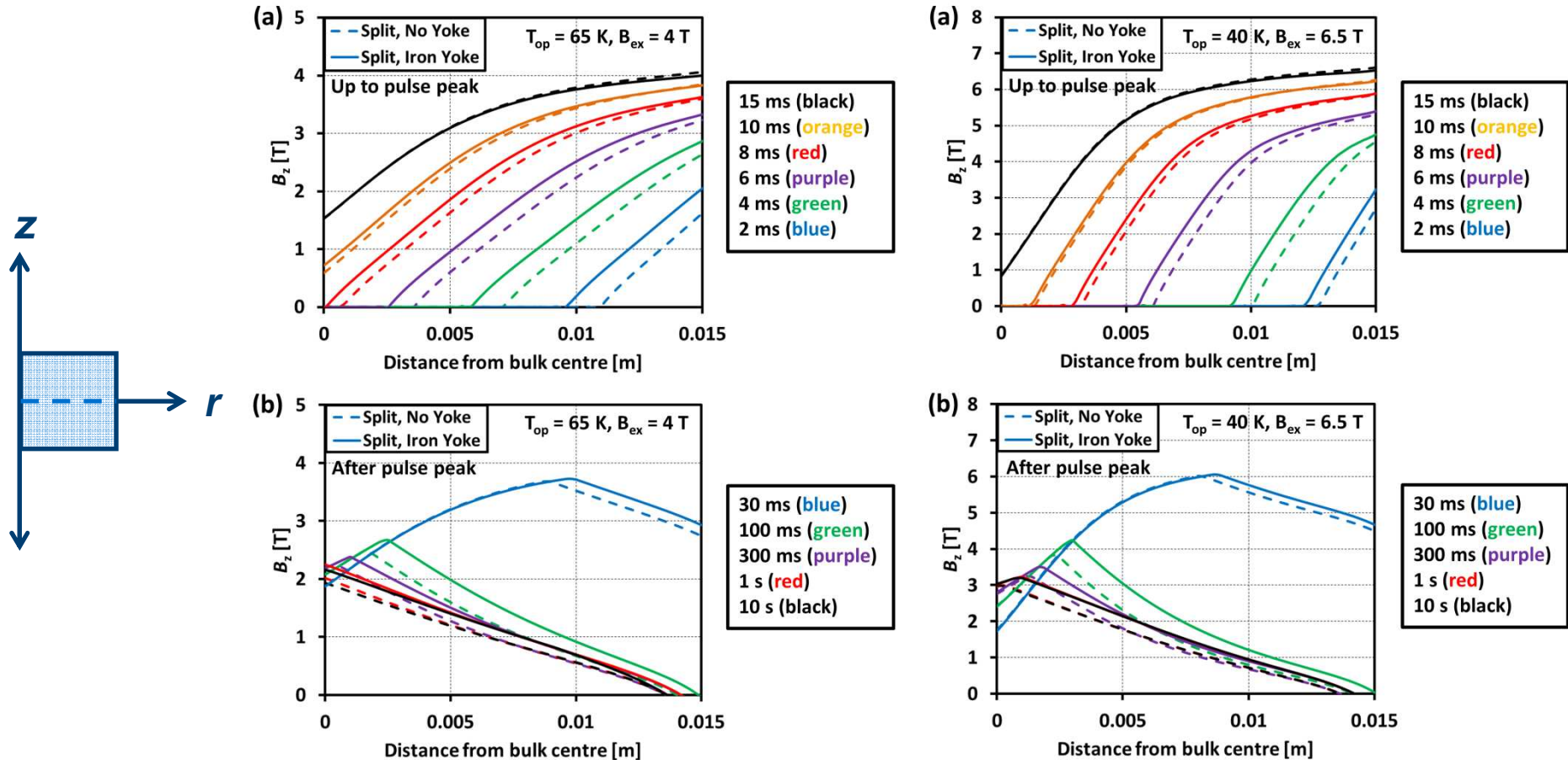
Trapped field & T_{ave} with time, incl. yoke removed



Ainslie, Fujishiro et al. *Supercond. Sci. Technol.* **29** (2016) 074003

Split Coil PFM with an Iron Yoke

Magnetic flux entry & exit during & after pulse – split coil with & without iron yoke



Ainslie, Fujishiro et al. *Supercond. Sci. Technol.* **29** (2016) 074003

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Thank you for listening

Contact email:

mark.ainslie@eng.cam.ac.uk

Website:

<http://www.eng.cam.ac.uk/~mda36/>