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Superconductors (HTS MODELLING 2016)

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Modelling pancake coils and solenoid coils wound
with coated conductors for electromagnetic field
analyses: using exact coil geometry or using
axisymmetric approximation

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This work was supported in part by Japan Science and Technology Agency under Strategic
Promotion of Innovative Research and Development Program (S-Innovation Program).



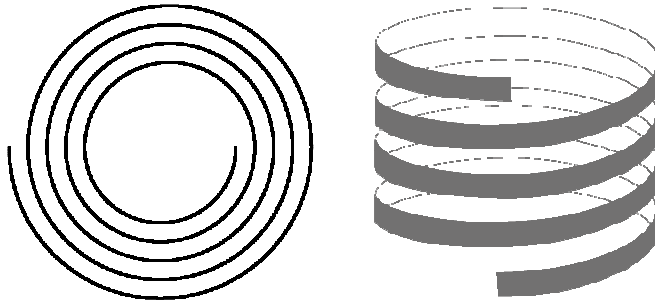
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Motivation

- Non-uniform current distributions in HTS tapes affect the field qualities of magnets wound with HTS tapes: shielding (screening) current issue.
- Electromagnetic field analyses are the popular approach to this issue.

Motivation (continued)

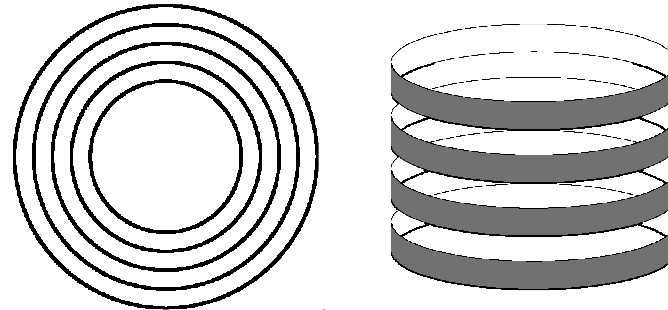
Exact geometry



Pancake coil: plane spiral
Solenoid coil: helix

A concern has been raised over the axisymmetric approximation for shielding current calculation (by Ueda et al.).

Axisymmetric approximation

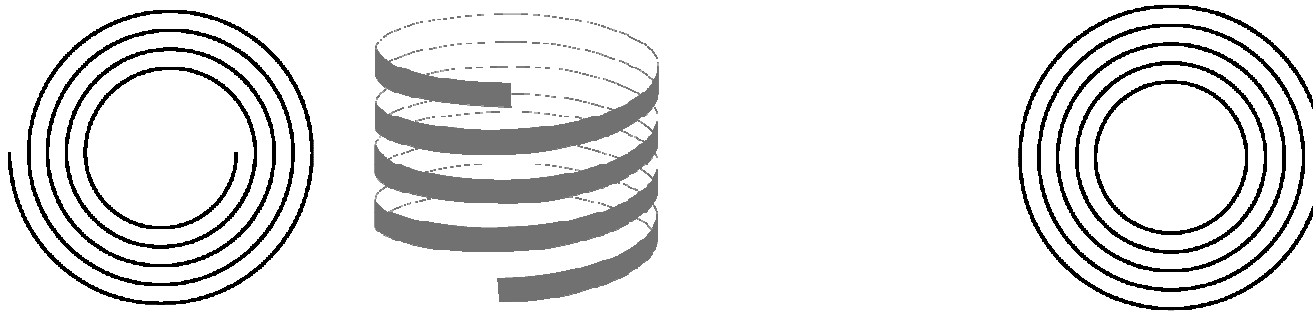


Pancake coil: nested one-turn loops
of coated conductors
Solenoid coil: stacked one-turn loops
of coated conductors

- Cross-sectional analyses
- Reduction of CPU time and memory

Objective

Comparison between the model based on the exact coil geometry and the model based on the axisymmetric approximation



- ✓ Focus on shielding-current-induced field (SCIF)
- ✓ Detailed results on stacked-pancake coils
- ✓ Some results on a multilayered solenoid coil: the difference from the pancake coils

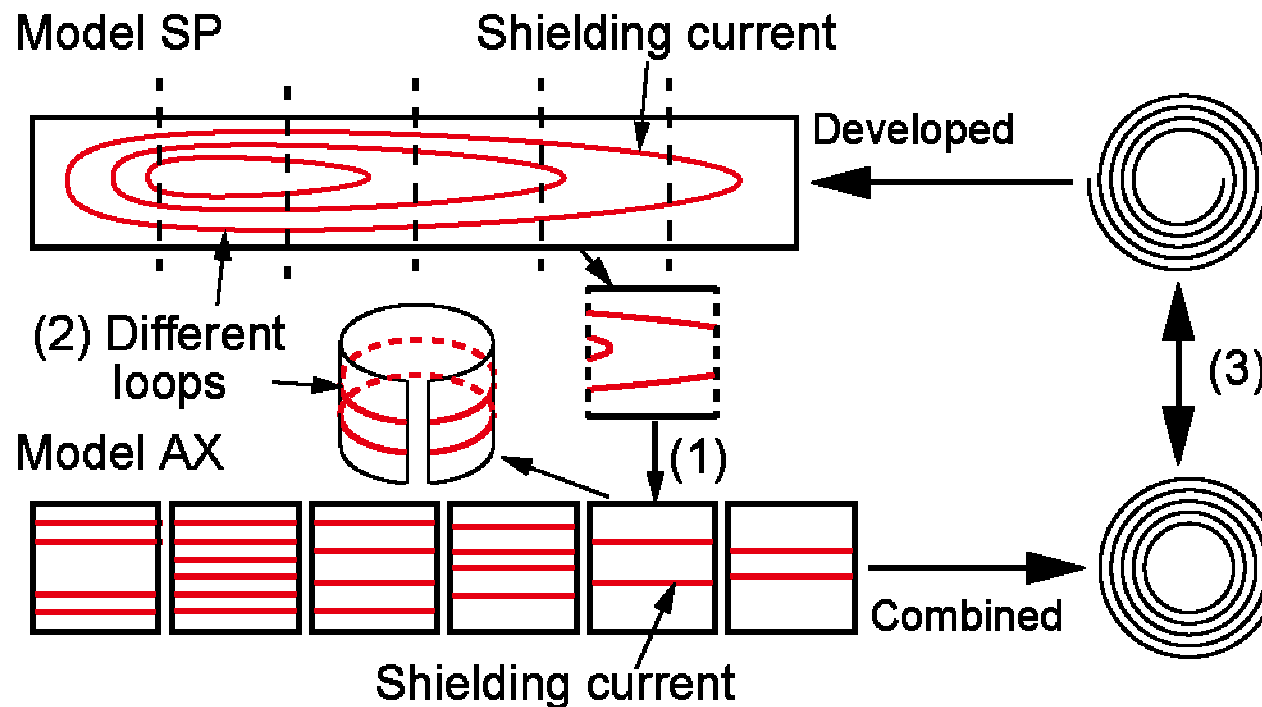
Outline

- Stacked pancake coils
 - Key issues of the concern on the axisymmetric approximation
 - Models for analyses
 - Results
- Multilayered solenoid coil
 - Model for analyses
 - Results
- Conclusion

Pancake coils

Key issues of the concern on the axisymmetric approximation

Key issues of concern on axisymmetric approximation



1. Different current path: neglecting transverse $J \rightarrow$ SCIF
2. Different inductance $\rightarrow$ decay time constant of shielding current
3. Different tape geometry (even if assuming uniform current distribution)

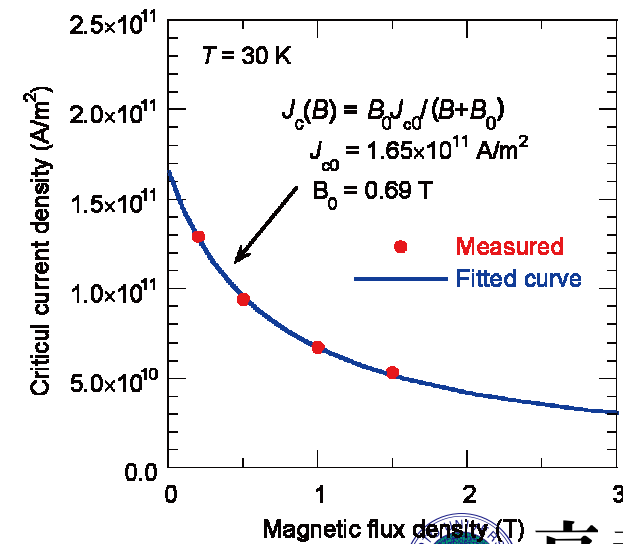
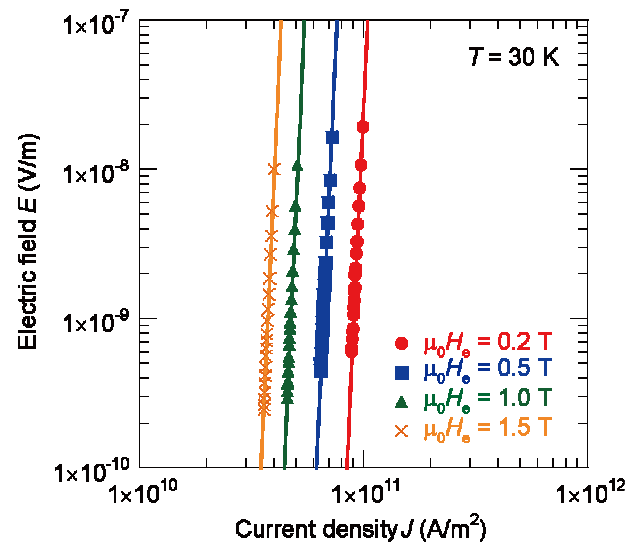
Models for analyses



E – J and J_c – B_n characteristics

$$E = E_0 \left(\frac{J}{J_c} \right)^n \quad J_c = J_{c0} \frac{B_0}{B_0 + |B_n|}$$

Parameters	
E_0	1×10^{-4} V/m
n	32.6
J_{c0}	1.65×10^{11} A/m ²
B_0	0.69 T

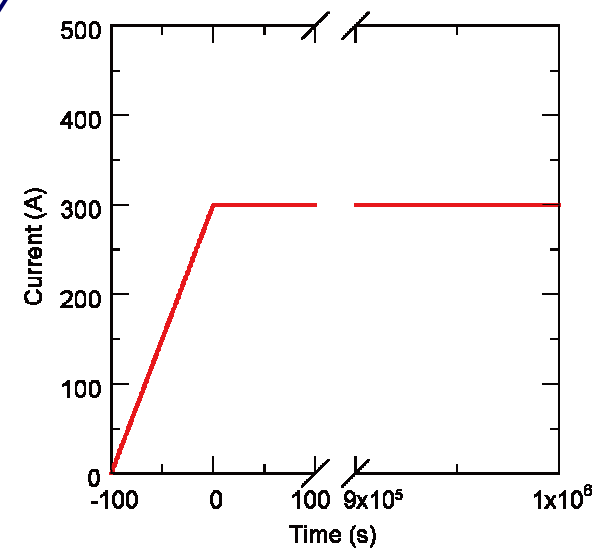
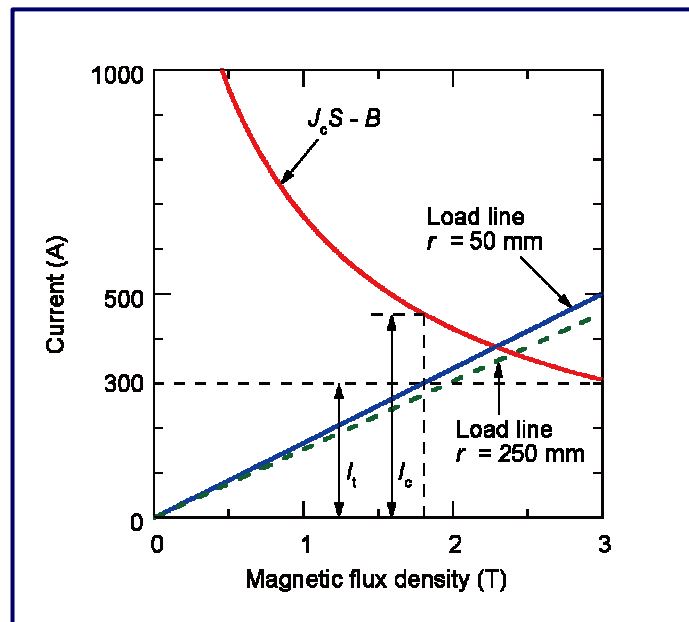


Stacked pancake coils for analyses

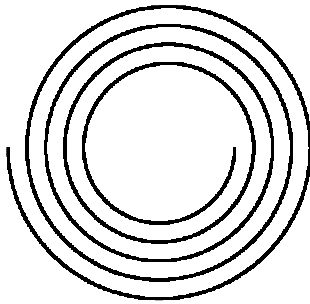
	PC-S	PC-L
Conductor width	5 mm	
Conductor / SC layer thickness	0.2 mm / 2 μ m	
Inner / outer radius	50 mm / 62.5 mm	250 mm / 262.5 mm
Number of turn per PC	50	50
Turn separation	0.05 mm	0.05 mm
Conductor length per PC	18 m	80.4 m
Number of PCs	30	30
Separation between PCs	0.5 mm	
Magnetic field at coil center	2.83 T @ 300 A	1.05 T @ 300 A
Maximum normal field component	1.8 T @ 300 A	1.97 T @ 300 A
Load ratio of coil @ 300 A	0.66	0.70

Operating point on $I - B$ plots and excitation pattern

	PC-S	PC-L
Magnetic field at coil center	2.83 T @ 300 A	1.05 T @ 300 A
Maximum normal field component	1.8 T @ 300 A	1.97 T @ 300 A
Load ratio of coil @ 300 A	0.66	0.70

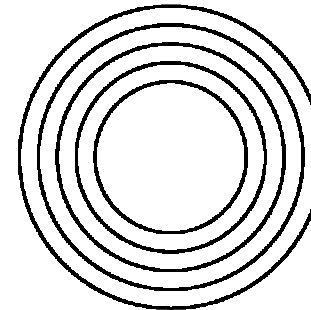


Two models for analyses



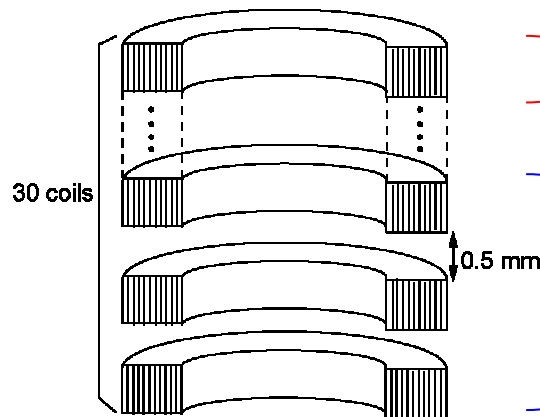
Model SP

(Exact spiral geometry)



Model AX

(Nested axisymmetric one-turn loops)



Analyzing one PC

Other PCs are represented by line currents in order to apply the field to the analyzed PC

Pancake by
pancake analysis

Repeated for
all PCs

Governing equation and constitutive equation

〔Faraday's law〕

$$\nabla \times \mathbf{E} + \frac{\partial \mathbf{B}}{\partial t} = 0$$

〔Biot-Savart's law〕

$$\mathbf{B} = \frac{\mu_0}{4\pi} \int_V \frac{\mathbf{J} \times \mathbf{r}}{r^3} dV$$

〔Definition of current vector potential〕

$$\mathbf{J} = \nabla \times \mathbf{T}$$

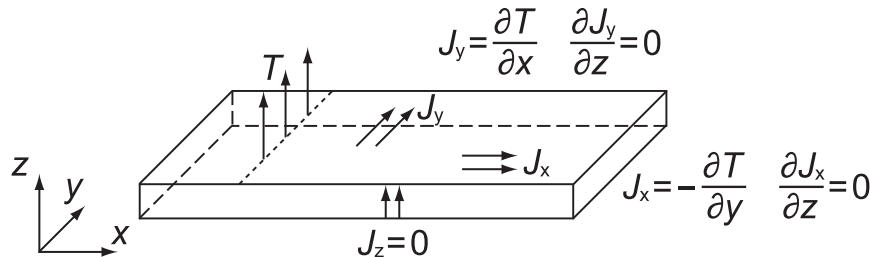
〔Extended Ohm's law〕

$$\mathbf{E} = \mathbf{J} / \sigma(\mathbf{J})$$

Equivalent conductivity derived from E-J characteristic: Ex. $E = E_0 (J / J_c)^n$

Thin strip approximation

High cross-sectional aspect ratio of coated-conductor allows its use.



$$\nabla \times \left(\frac{1}{\sigma} \nabla \times \mathbf{T}_f \right) + \frac{\mu_0}{4\pi} \frac{\partial}{\partial t} \int_V \frac{(\nabla \times \mathbf{T}_s) \times \mathbf{r}}{r^3} dV = 0$$

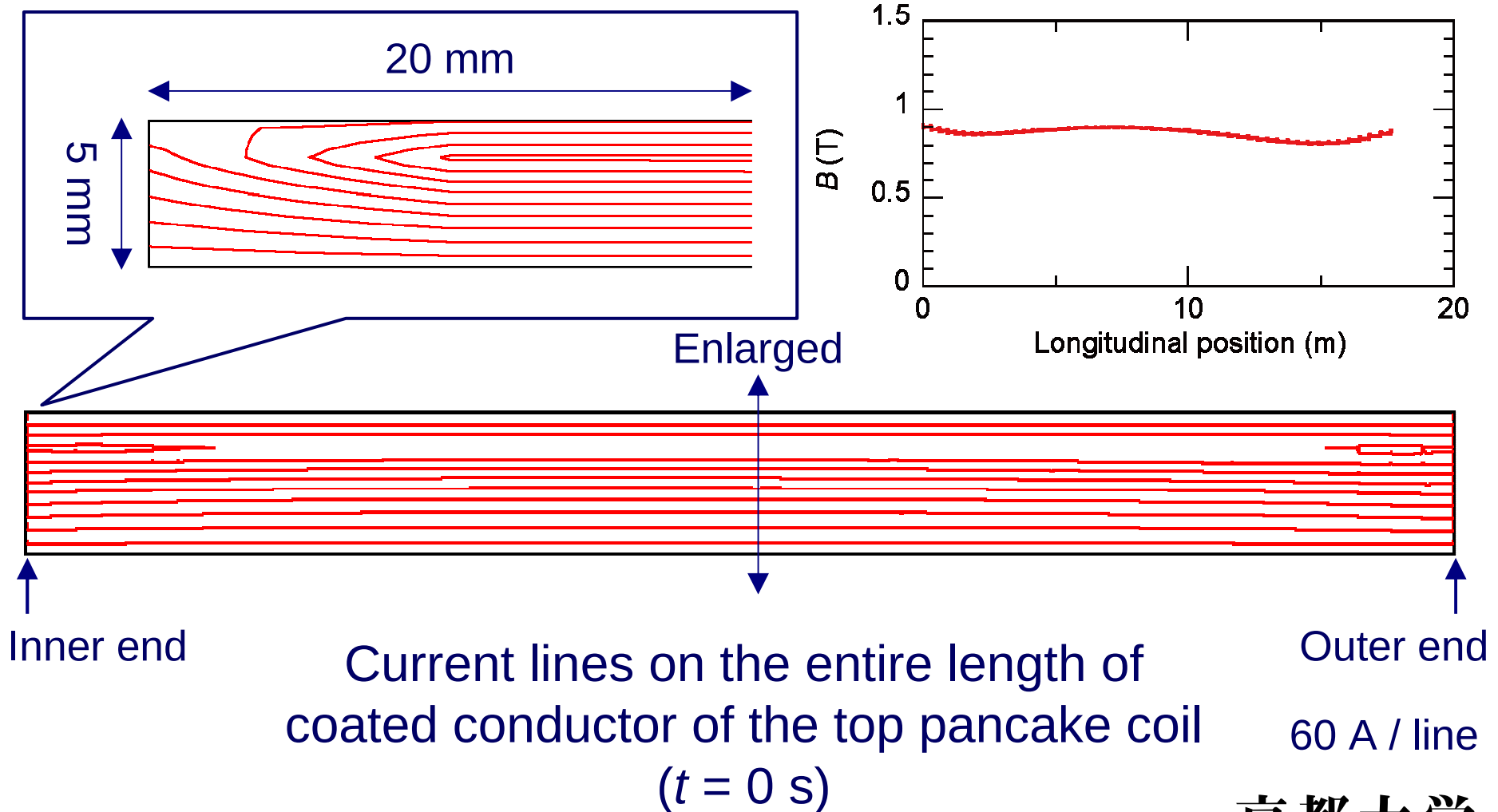
Integrate along thickness of coated-conductor

$$\nabla \times \left(\frac{1}{\sigma} \nabla \times \mathbf{T}_f \right) + \frac{\mu_0 t_s}{4\pi} \frac{\partial}{\partial t} \int_S \frac{(\nabla \times \mathbf{T}_s) \times \mathbf{r}}{r^3} dS = 0$$

Results

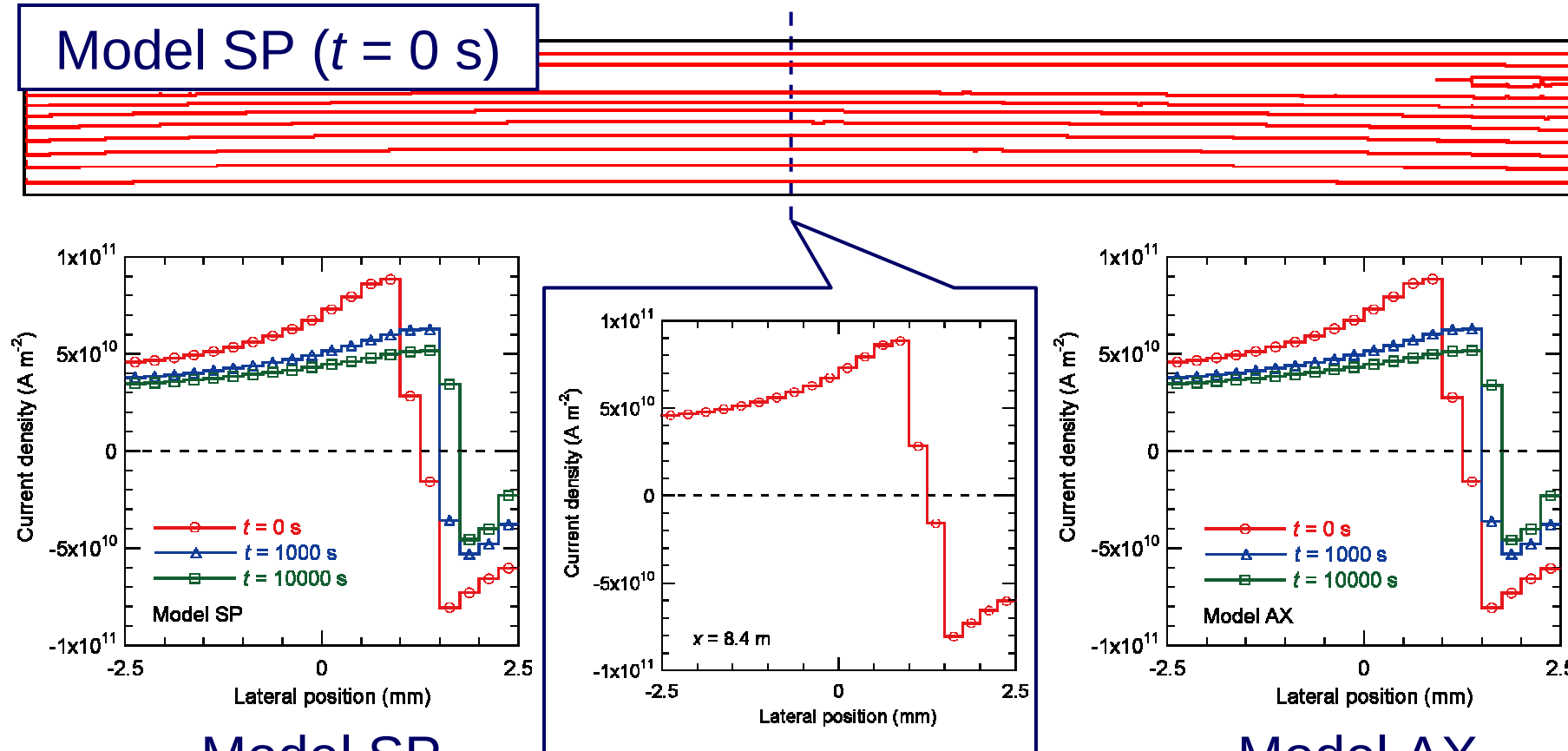


Model SP: normal field component and current profile with induced shielding current



Lateral current distributions (top pancake)

Model SP ($t = 0$ s)



Model SP

Center of the 25th turn
of top pancake

Model AX

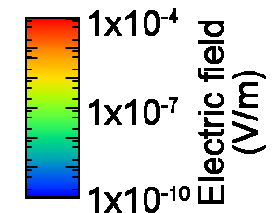
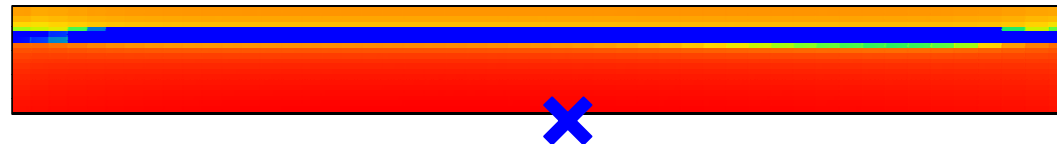
$$\Delta J/J \sim 10^{-4}$$



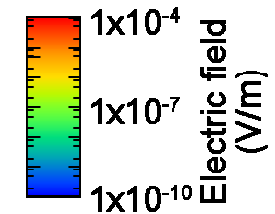
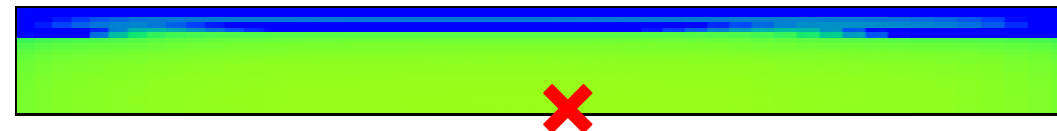
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Electric field and change in current distribution

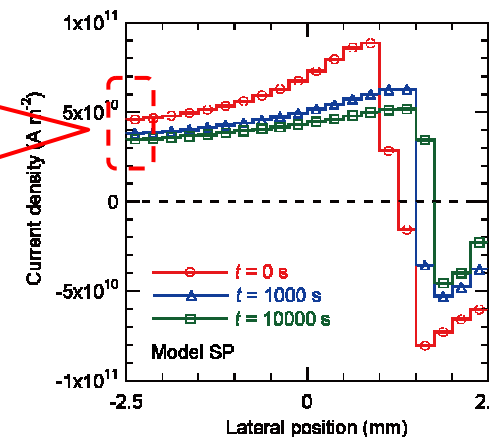
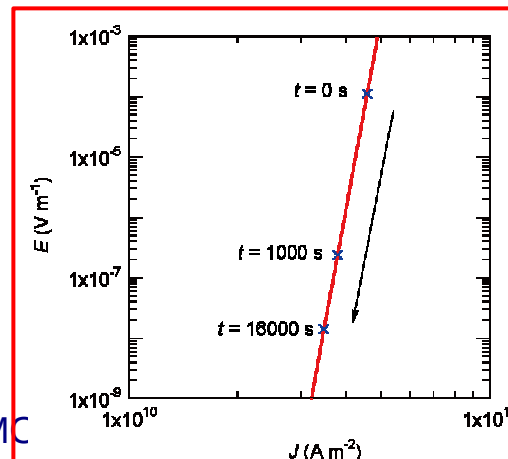
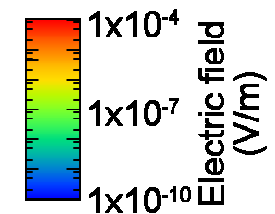
$t = 0 \text{ s}$



$t = 1000 \text{ s}$



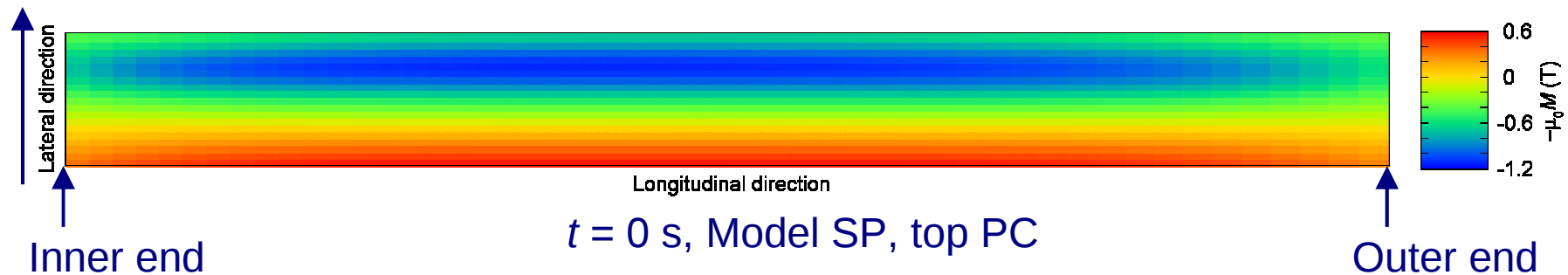
$t = 16000 \text{ s}$



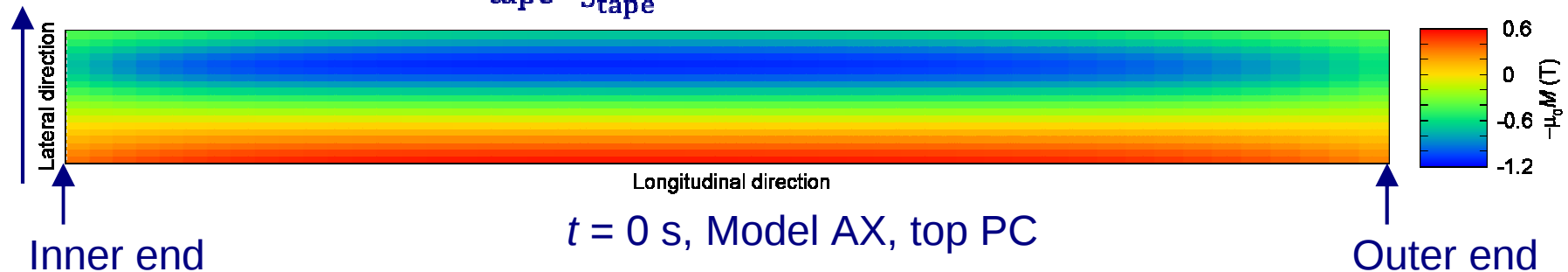
Magnetization distribution on coated conductor

$$-\mu_0 M = B - \mu_0 H$$

$-\mu_0 M$: magnetization
 B : Magnetic flux density
 $\mu_0 H$: Field generated by other PCs

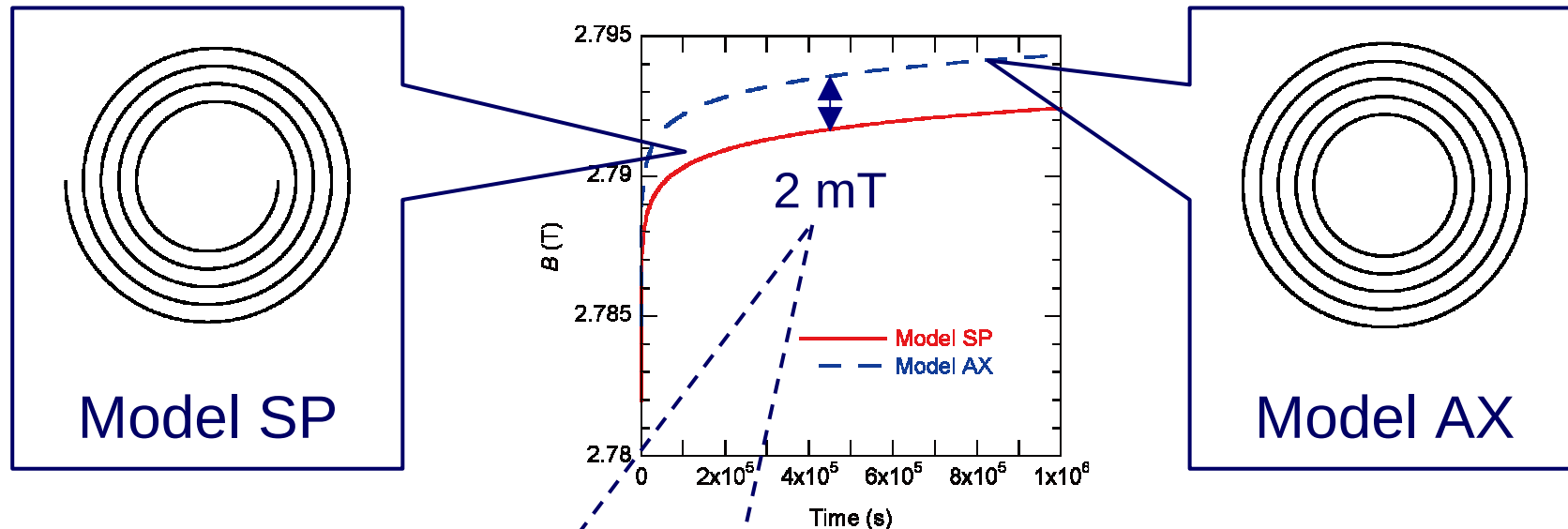


$$-\mu_0 M_{\text{tape}} = \frac{1}{S_{\text{tape}}} \int_{S_{\text{tape}}} -\mu_0 M dS = 345.9 \text{ mT}$$



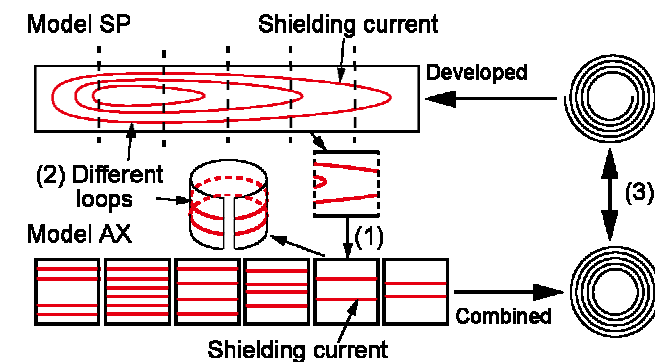
$$-\mu_0 M_{\text{tape}} = \frac{1}{S_{\text{tape}}} \int_{S_{\text{tape}}} -\mu_0 M dS = 346.3 \text{ mT}$$

Temporal evolution of magnetic field at coil center



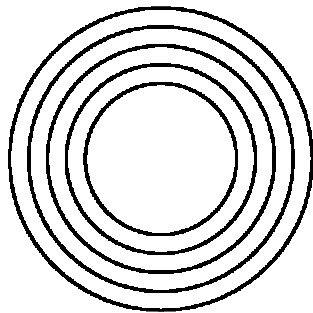
Is this difference caused by

- different current path: neglecting transverse J ?
- different geometry?



Compensation of influence of different geometry

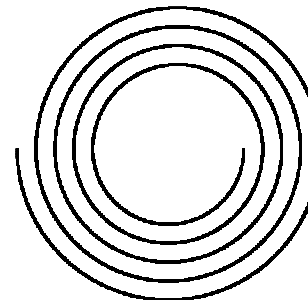
Model AX



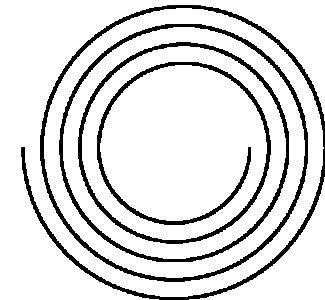
2. Projecting the current distributions on the spiral coated conductor



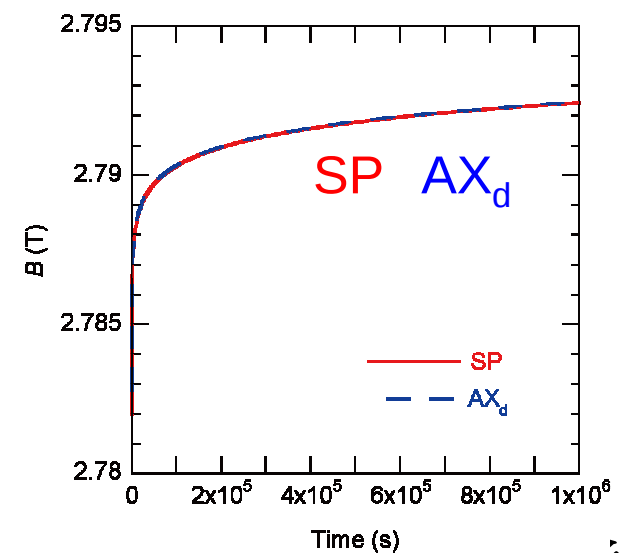
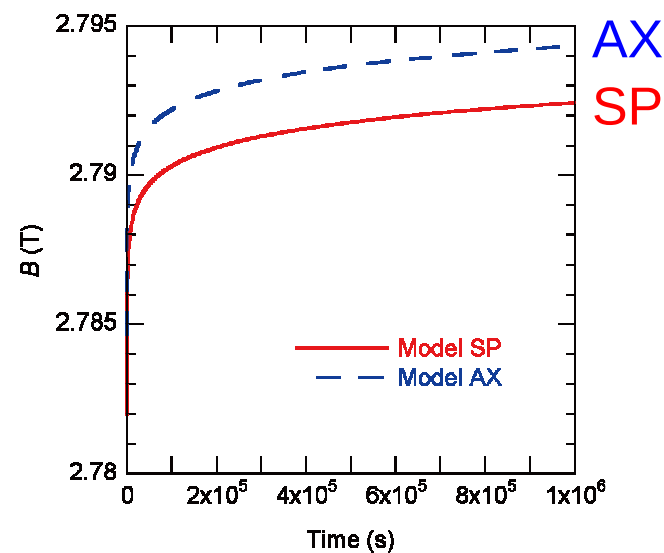
AX_d



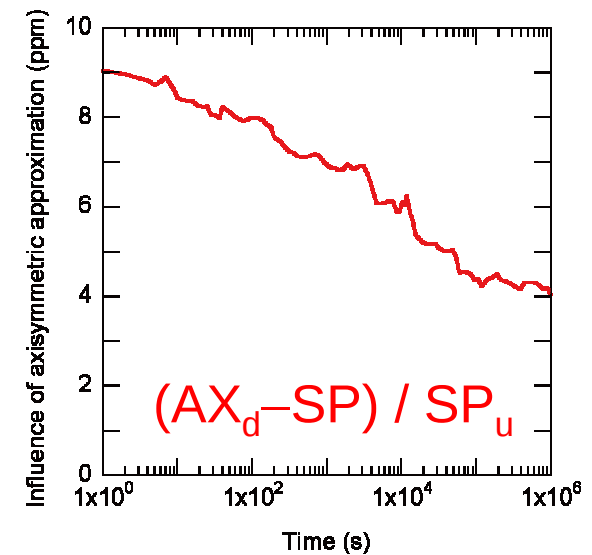
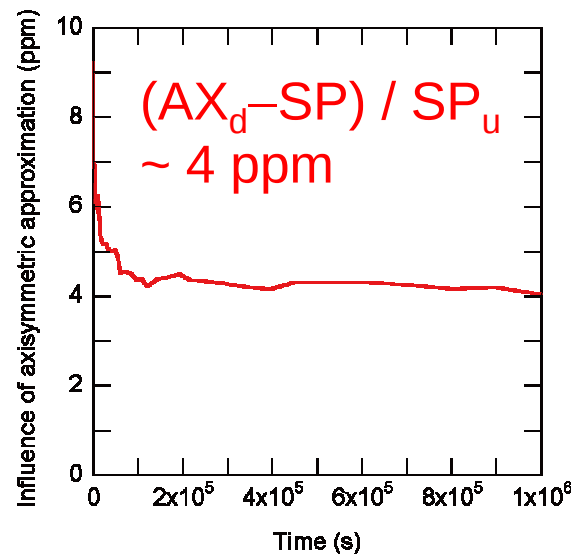
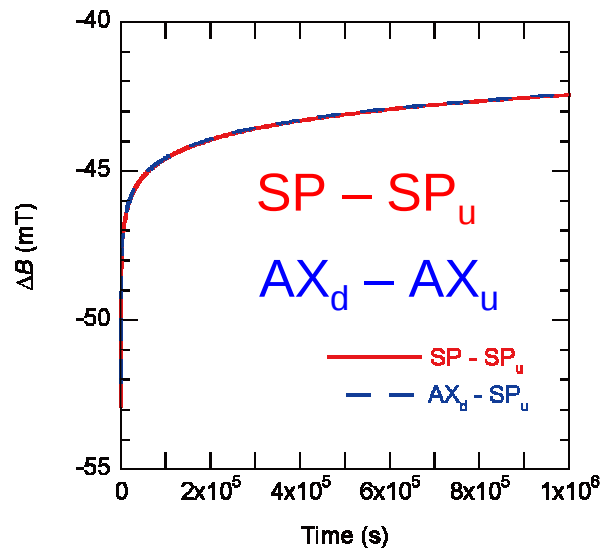
Model SP



1. Calculating current distributions by using model AX



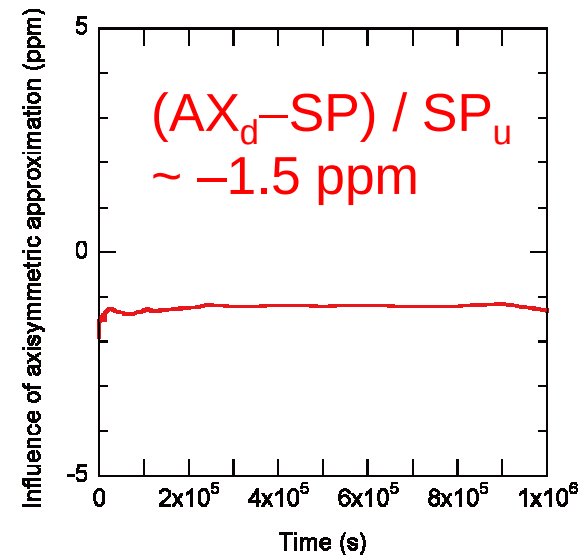
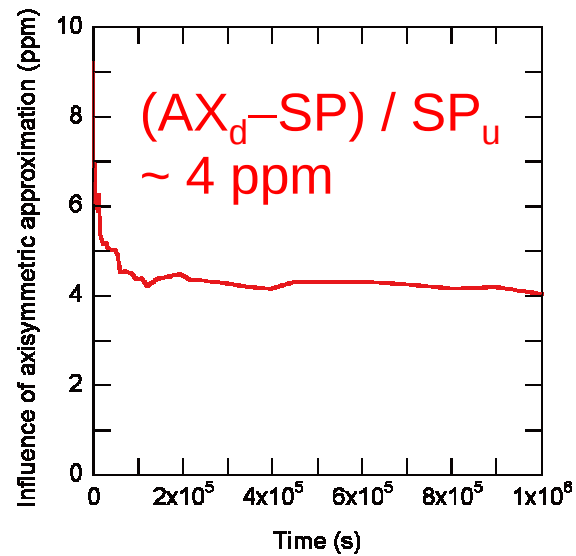
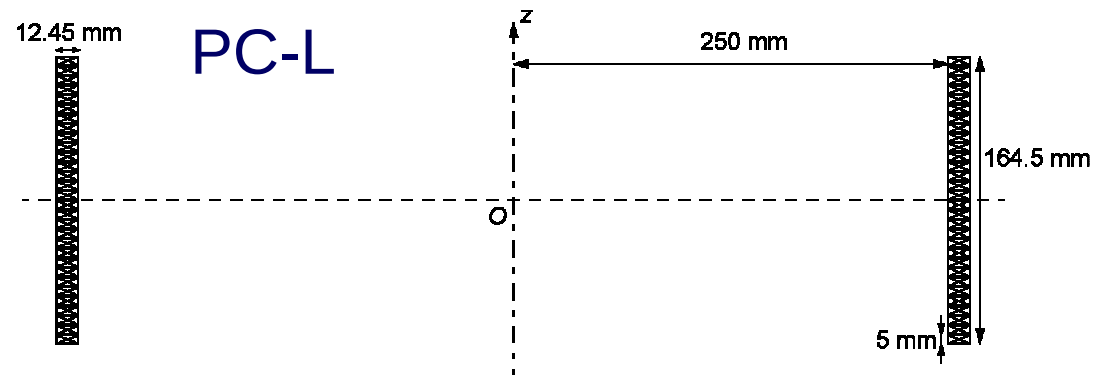
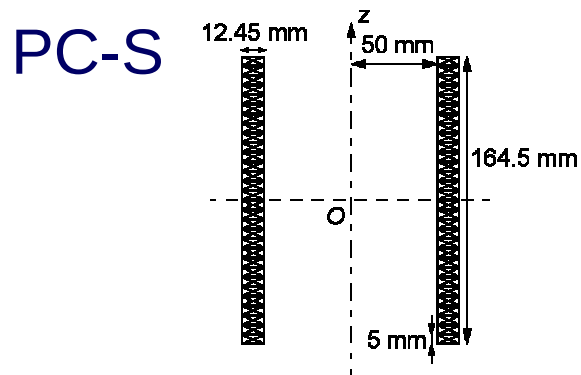
Influence of SCIF



Influence of SCIF

Difference between AX_d and SP
normalized by the field calculated
with uniform current distribution

Comparison between PC-S and PC-L



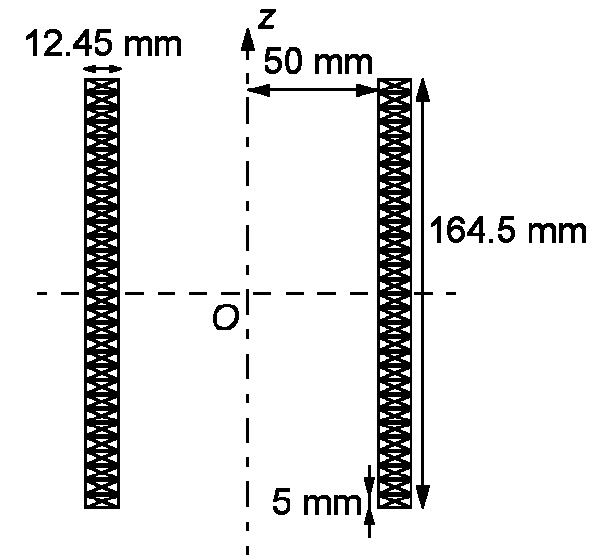
Solenoid coil

Models for analyses

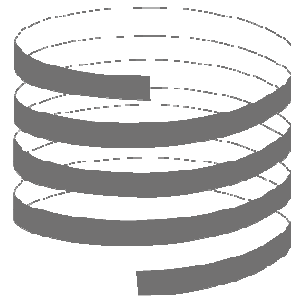


Multilayered solenoid coil

	SL
Conductor width	5 mm
Conductor / SC layer thickness	0.2 mm / 2 μ m
Inner / outer radius	50 mm / 62.5 mm
Number of turn per layer	30
Turn separation	0.5 mm
Conductor length per layer	12 m
Number of layers	50
Separation between layer	0.05 mm



Models for analyses



Model HX (Exact spiral geometry)

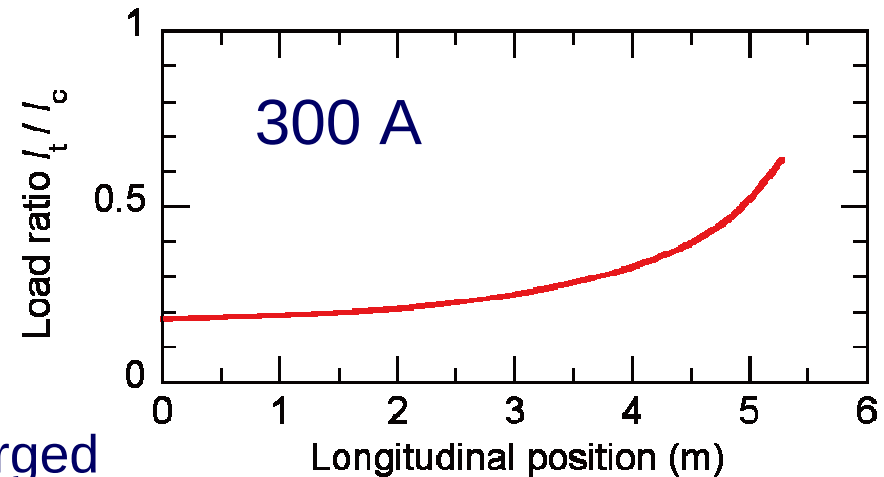
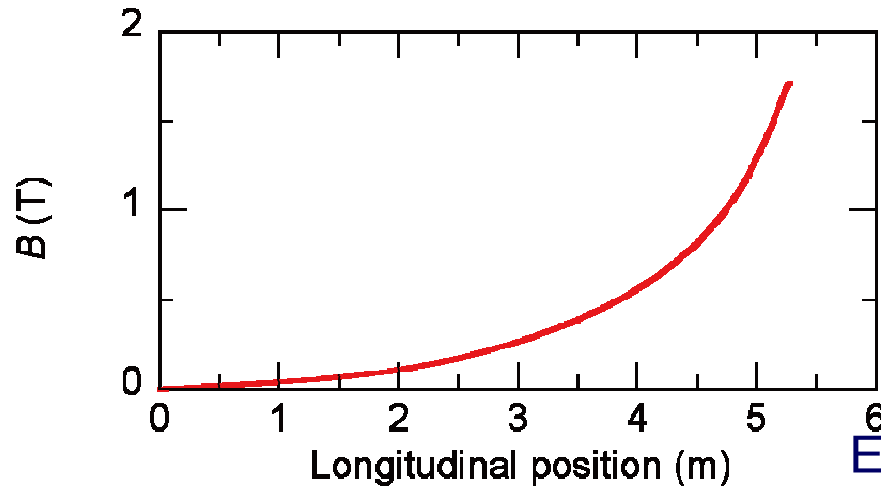
- ✓ The layer-by-layer analysis applied: not analyzed layers represented by sets of line currents in order to apply the field to the analyzed layer

Results

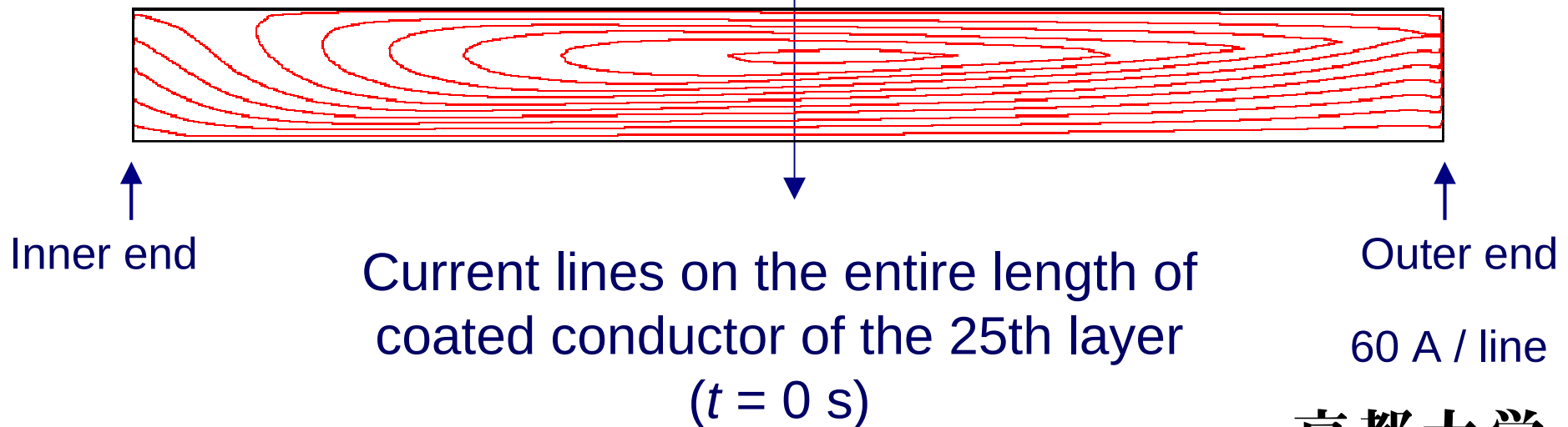


Model HX: normal field component, load ratio and current profile with induced shielding current

SL



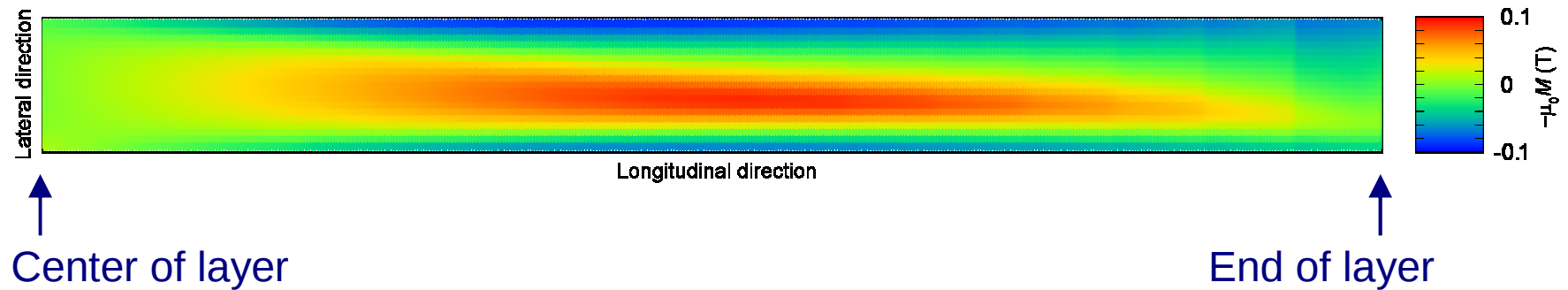
Enlarged



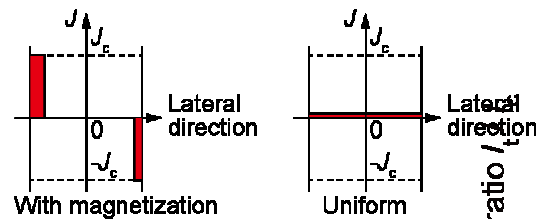
Magnetization distribution on coated conductor

$$-\mu_0 M = B - \mu_0 H$$

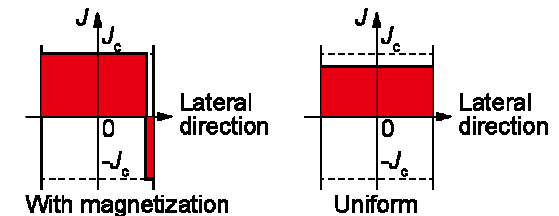
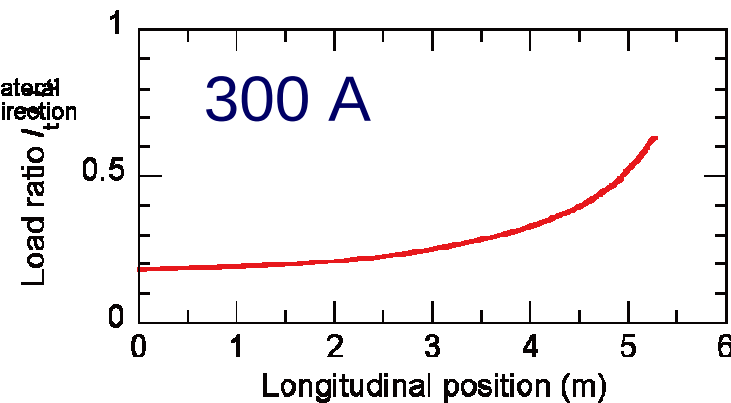
$-\mu_0 M$: magnetization
 B : Magnetic flux density
 $\mu_0 H$: Field generated by other PCs



$t = 0$ s, Model HX, the 25th layer



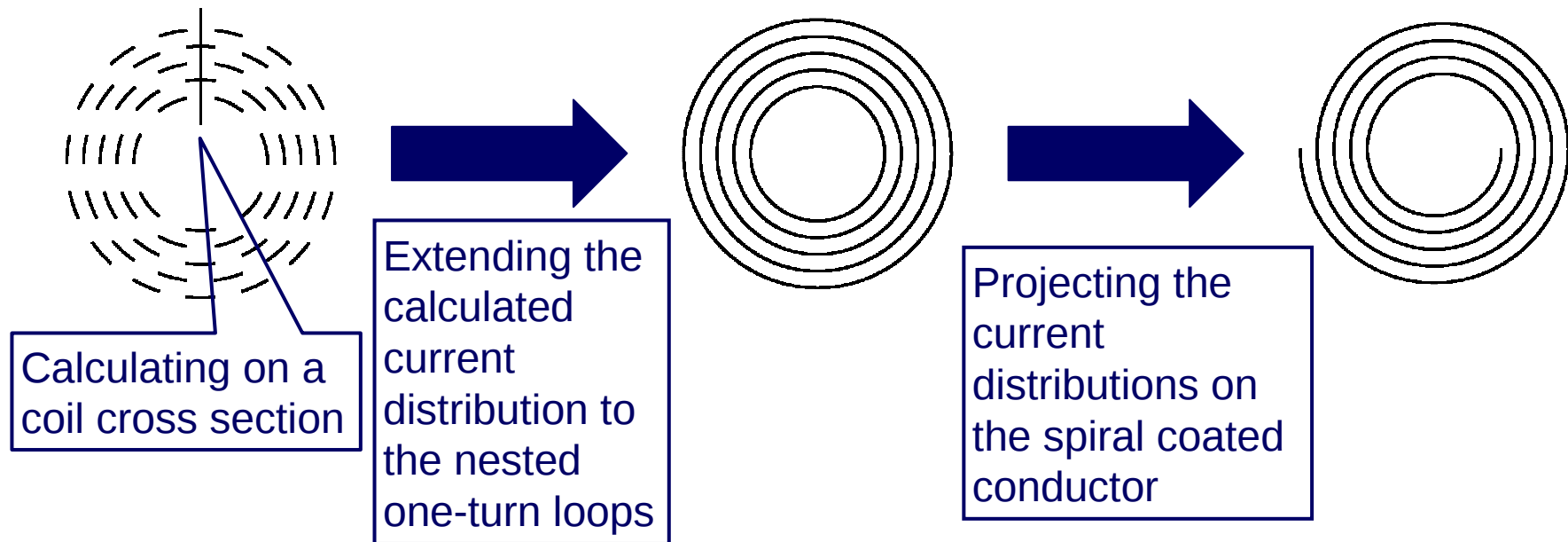
Low load ratio



High load ratio

Conclusion

- The following scheme enables us to calculate the magnetic field of a stacked pancake coils accurately and with less CPU time and memory.

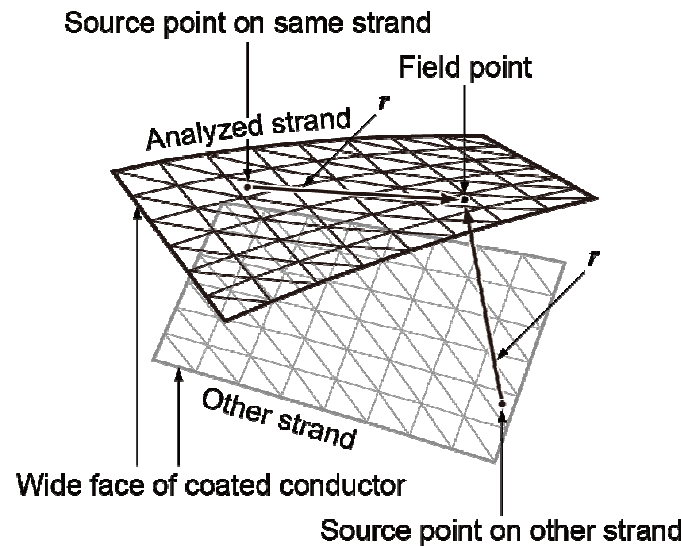


- The approximation of nested one-turn loops is also useful for non-axisymmetric coils such as saddle-shape coils.

Back-up slides



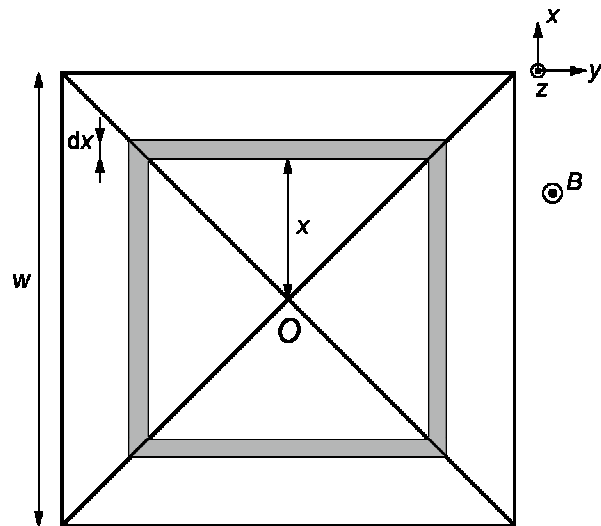
Consideration of three-dimensionally-curved coated conductors



$$\nabla \times \left(\frac{1}{\sigma} \nabla \times \mathbf{n} T \right) \cdot \mathbf{n} + \frac{\partial}{\partial t} \left(\frac{\mu_0 t_s}{4\pi} \int_{S'} \frac{(\nabla \times \mathbf{n}' T') \times \mathbf{r} \cdot \mathbf{n}}{r^3} dS' + \mathbf{B}_{\text{ext}} \cdot \mathbf{n} \right) = 0$$

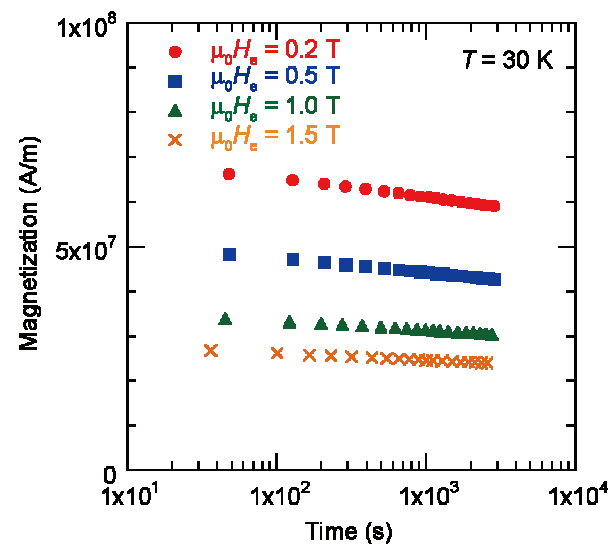
This term representing $\mathbf{B}$ in Faraday's law is calculated by Biot-Savart's law based on currents on arbitrary 3D-shaped conductors

Magnetization measurements to determine E – J

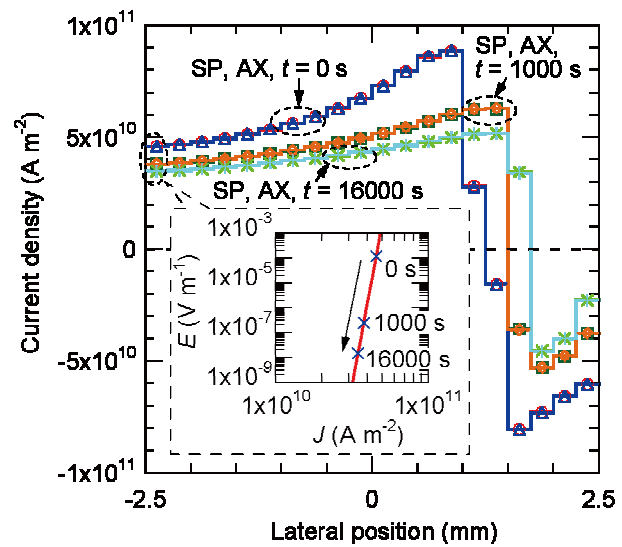


$$J = \frac{6m}{w^3 d} \quad E = -\frac{\mu_0 G}{4wd} \cdot \frac{dm}{dt}$$

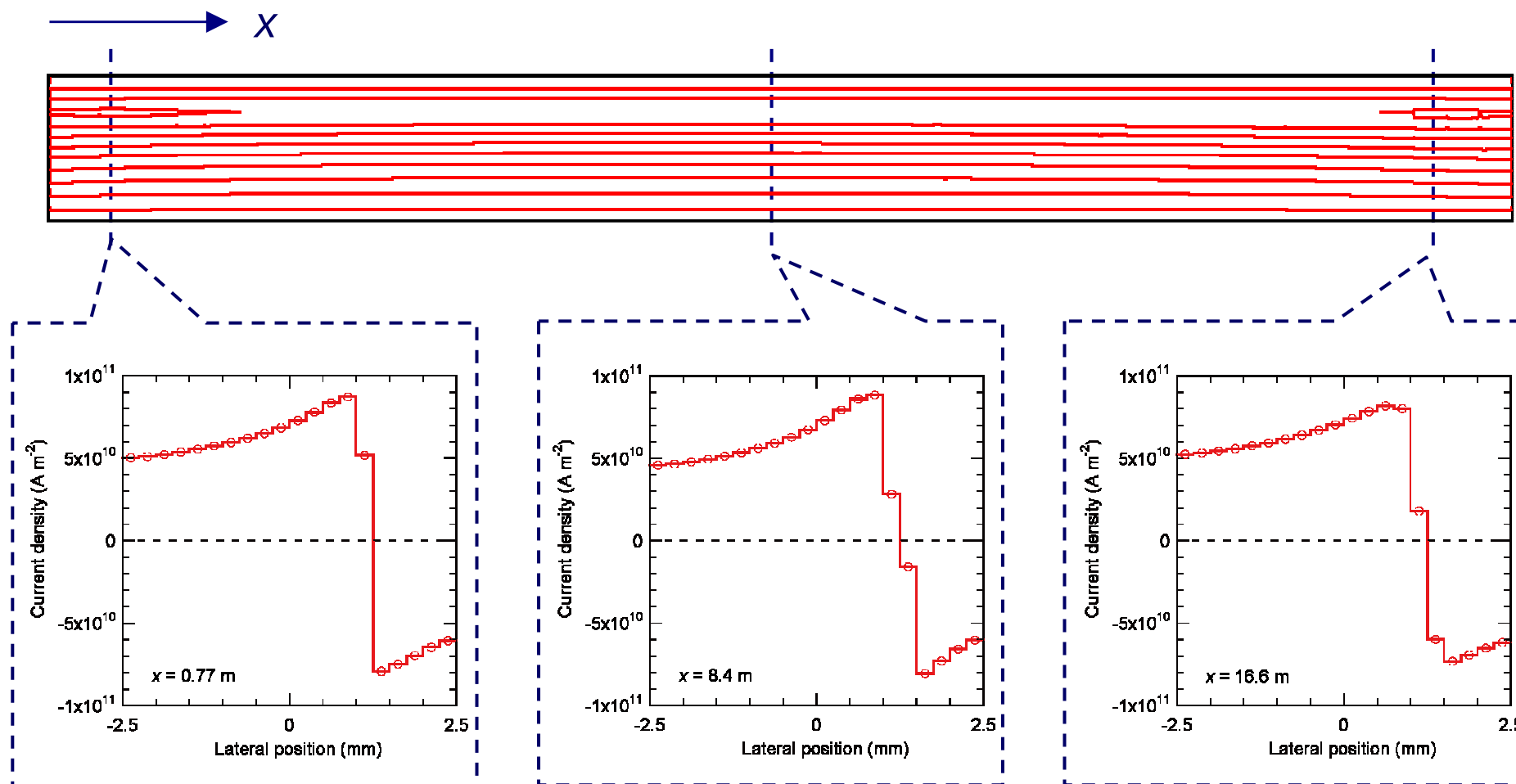
d : thickness of superconductor
 G : geometrical factor 7.17×10^{-4}
 m : magnetization of sample



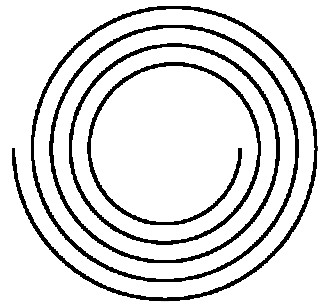
Figures not used



Lateral current distributions



Influence of different geometry (uniform current)



Model SP

**Magnetic field calculated with
uniform current distribution**

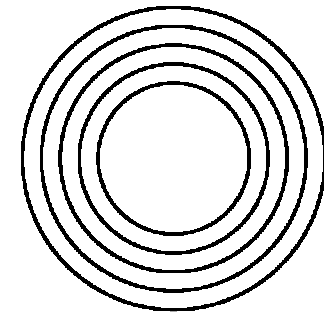
Model SP (SP_u)

2.83488 T

Model AX (AX_u)

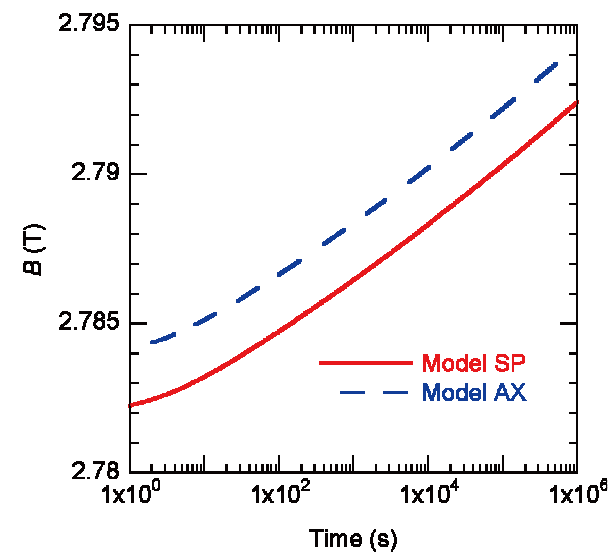
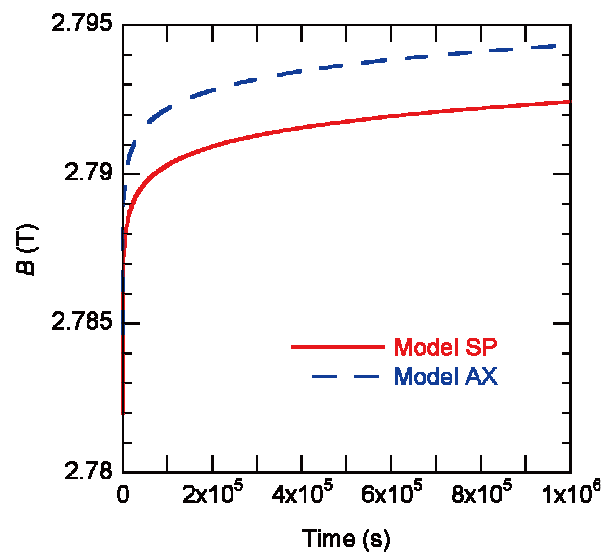
2.83686 T

Difference ~ 2 mT



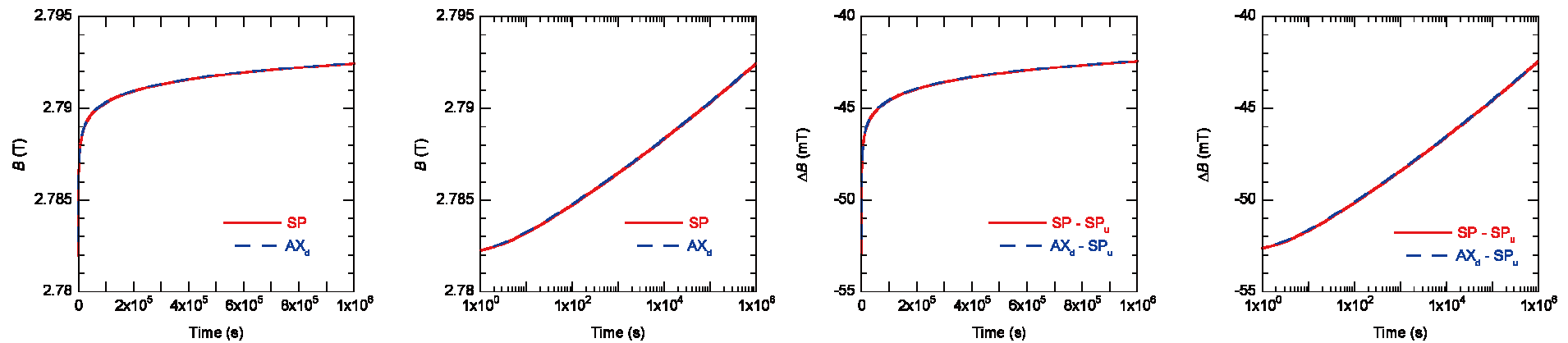
Model AX

Temporal evolution of magnetic field at coil center



Log plot

SP, AX, AX_d (linear, log, values)



遮蔽電流の影響を考慮した
場合の発生磁場 (電磁界解析結果)

Model SP (SP)	Model AX (AX)	Model AX (AX _d)
2.78198 T ($t = 0$ s)	2.78389 T ($t = 0$ s)	2.78201 T ($t = 0$ s)

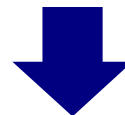
軸対称近似のコイル発生磁場に対する影響
(AX_d-SP) / SP_u = 9.2 ppm

3. 解析結果

異なる時刻での軸対称近似が磁場に与える影響

遮蔽電流の影響を考慮した 場合の発生磁場 (電磁界解析結果)			
ランプアップ終了 からの経過時間 t	SP	AX_d	$(AX_d - SP) / SP_u$
0 s	2.7820 T	2.7820 T	9.2 ppm
100 s	2.7847 T	2.7848 T	8.0 ppm
10000 s	2.7883 T	2.7884 T	6.1 ppm
1000000 s	2.7924 T	2.7924 T	4.0 ppm

軸対称近似のコイル発生磁場に対する影響の時間変化は非常に小さい

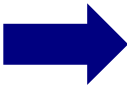


軸対称近似の適用による遮蔽電流の減衰時定数への影響は小さい

3. 解析結果

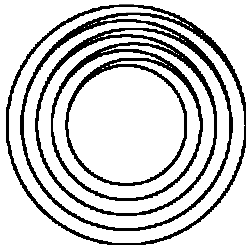
PC群中心での磁場のモデル間での差異に関する検討

線材内に電流が一様に分布していた 場合の発生磁場	
Model SP (SP_u)	Model AX (AX_u)
1.05396 T	1.05442 T



コイル形状が異なることによって
0.5 mTの差が発生する。
(半径が大きくなった分、コイル形状
の違いの影響も小さくなった)

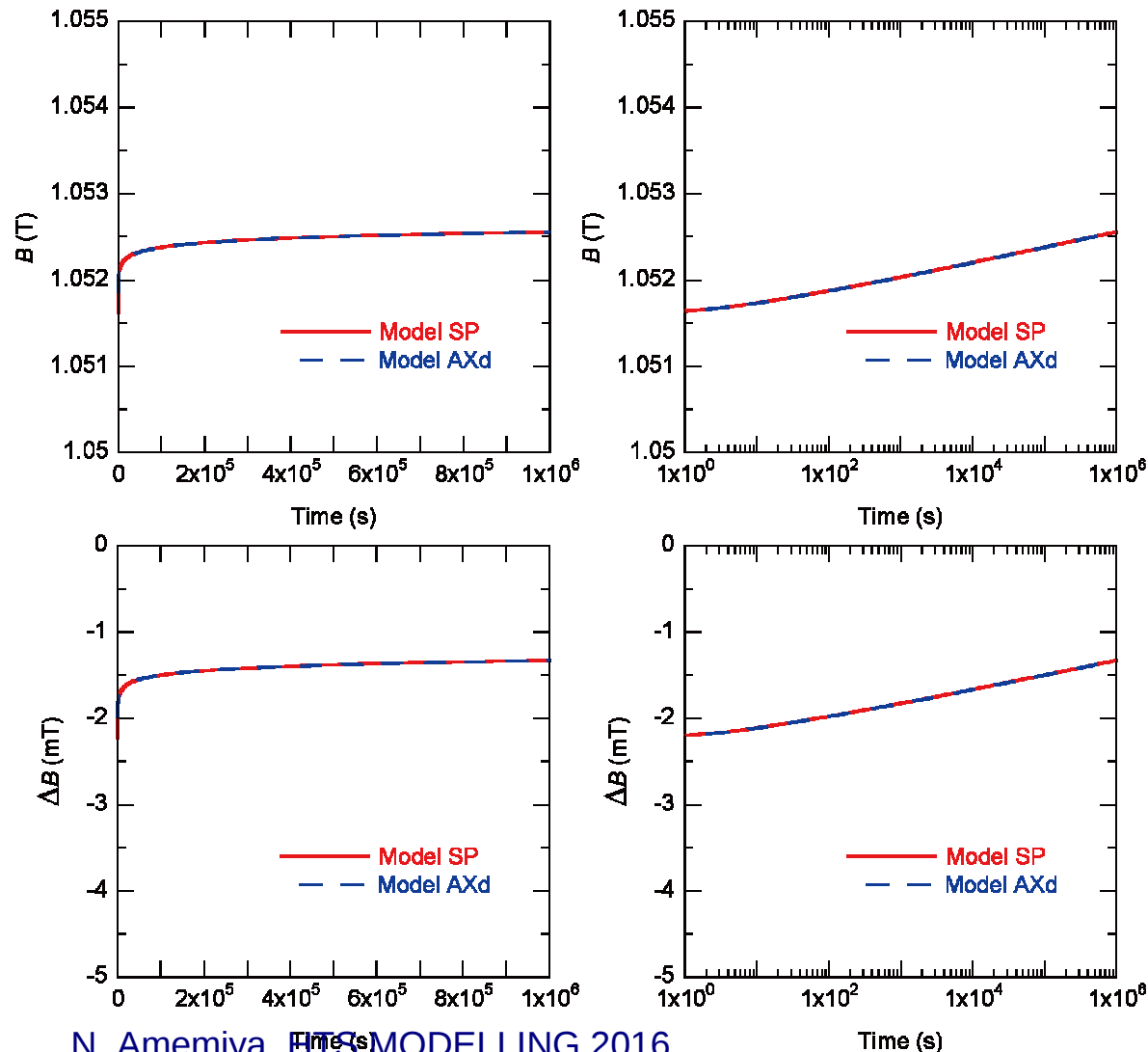
遮蔽電流の影響を考慮した 場合の発生磁場 (電磁界解析結果)		
Model SP (SP)	Model AX (AX)	Model AX (AX_d)
1.05162 T ($t = 0$ s)	1.05208 T ($t = 0$ s)	1.05161 T ($t = 0$ s)



軸対称近似のコイル発生磁場に対する影響
 $(AX_d - SP) / SP_u = -1.9$ ppm

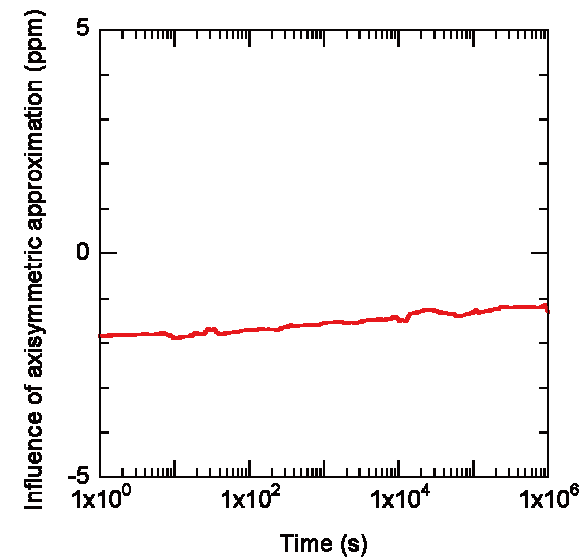
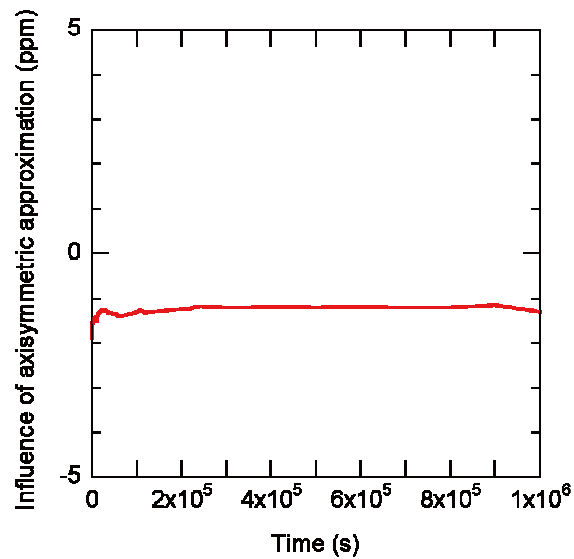
軸対称近似厳密形状得基ねば線群内電流写像を

遮蔽電流の磁場への影響の時間変化



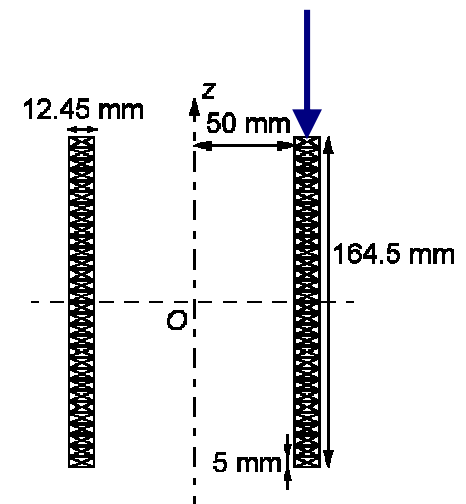
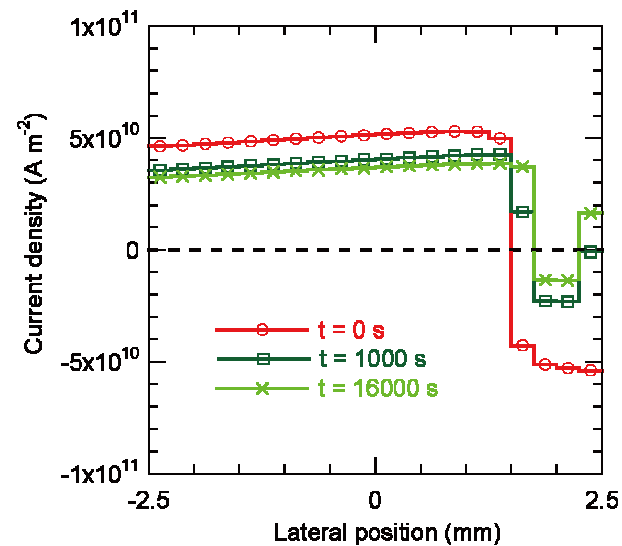
$r = 50$ mmの場合よりも
遮蔽電流磁場の減衰は小さい

軸対称近似の影響の変化



コイル半径が5倍になったとき、軸対称近似の影響は約1/4倍になった。
→コイル径が大きくなればなるほど、軸対称近似による解析で十分になる可能性が高い

幅方向電流密度分布



ソレノイドコイル25層目、30ターン目中央

外部磁場の与え方が異なっているため直接比較できないが、パンケーキコイルの場合とは分布が多少異なる。
時間変化の様子はパンケーキコイルと同様

Lateral current distributions

