

TDGL Simulation on Dynamics of Helical Vortices in Thin Superconducting Wires in the Force-Free Configuration

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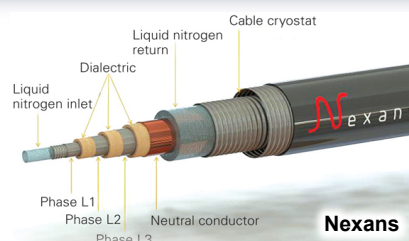
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1. Time-Dependent Ginzburg-Landau (TDGL) Equations
2. 2D Simulation by TDGL equations
3. Vortex Dynamics and Electromagnetic Response
4. Longitudinal Field Effect in the Force-Free Configuration
5. 3D Simulation: Transverse Configuration
6. 3D Simulation: Longitudinal (Force-Free) Configuration
7. Summary

Modelling of Superconductors

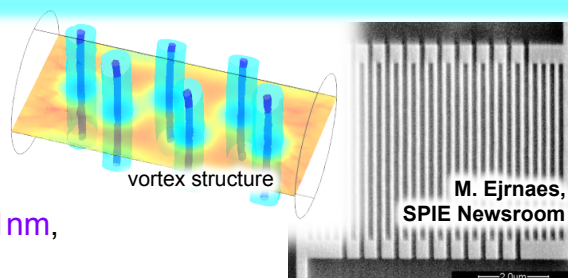
◆ Macroscopic Modelling (> mm)

- modelling of wires, conductors, and power devices
- critical state model, power-law model $E \sim J^n$



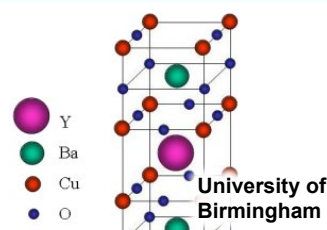
◆ Mesoscopic Modelling (nm ~ μm)

- modelling of nanoscale phenomena and electronics devices
- Ginzburg-Landau equation, London model, Josephson equation
- length scales: coherence length $\xi \sim 1\text{nm}$, penetration depth $\lambda \sim 0.1\mu\text{m}$



◆ Microscopic Theory (< nm)

- theory for electronic state and superconducting mechanism
- BCS theory, first-principle calculation



Time-Dependent Ginzburg-Landau Equation

A. Schmid, Phys. Kondens. Materie **5**, 302 (1966).

◆ Time-dependent Ginzburg-Landau (TDGL) equation:

- order parameter $\psi = |\psi| \exp(i\varphi)$; superfluid density $|\psi|^2$, phase φ
- ψ normalized ($|\psi| = 1$ for the equilibrium state at zero field)

$$\tau_0 \left(\frac{\partial}{\partial t} + i \frac{2\pi}{\phi_0} \Phi \right) \psi = \xi_0^2 \left(\nabla - i \frac{2\pi}{\phi_0} \mathbf{A} \right)^2 \psi + (1 - |\psi|^2) \psi$$

◆ Current density, $\text{rot rot } \mathbf{A} / \mu_0 = \mathbf{j}_s + \mathbf{j}_n$

$$\left\{ \begin{array}{l} \mathbf{j}_s = \frac{|\psi|^2}{\mu_0 \lambda_0^2} \left(\frac{\phi_0}{2\pi} \nabla \varphi - \mathbf{A} \right), \quad \rightarrow \text{fluxoid quantization} \\ \mathbf{j}_n = \sigma_n \mathbf{E} = -\sigma_n \left(\frac{\partial \mathbf{A}}{\partial t} + \nabla \Phi \right) \end{array} \right.$$

- upper critical field $B_{c2} = \frac{\phi_0}{2\pi \xi^2}$

- depairing current density

$$J_d = \frac{2\sqrt{2}}{3\sqrt{3}} \frac{B_c}{\mu_0 \lambda} = \frac{\phi_0}{3\sqrt{3} \pi \mu_0 \lambda^2 \xi}$$

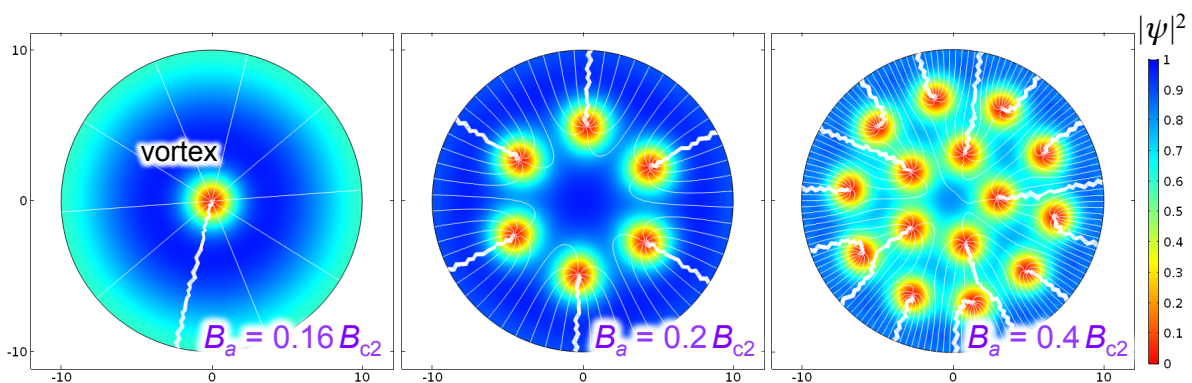
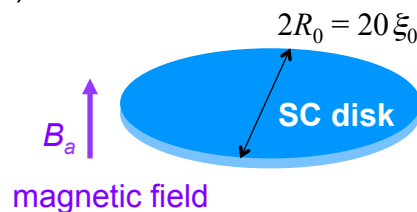
- τ_0 : GL relaxation time
- ξ_0 : coherence length
- λ_0 : penetration depth
- Φ : scalar potential
- \mathbf{A} : vector potential
- ϕ_0 : flux quantum

	NbTi	MgB ₂	YBa ₂ Cu ₃ O _y
T_c	9.5K	39K	95K
τ_0	0.3ps	0.08ps	0.03ps
ξ_0	4nm	5nm	0.4–3nm
λ_0	300nm	140nm	30–200nm

2D TDGL Simulation for SC Disks

◆ SC disk in a perpendicular magnetic field (2D)

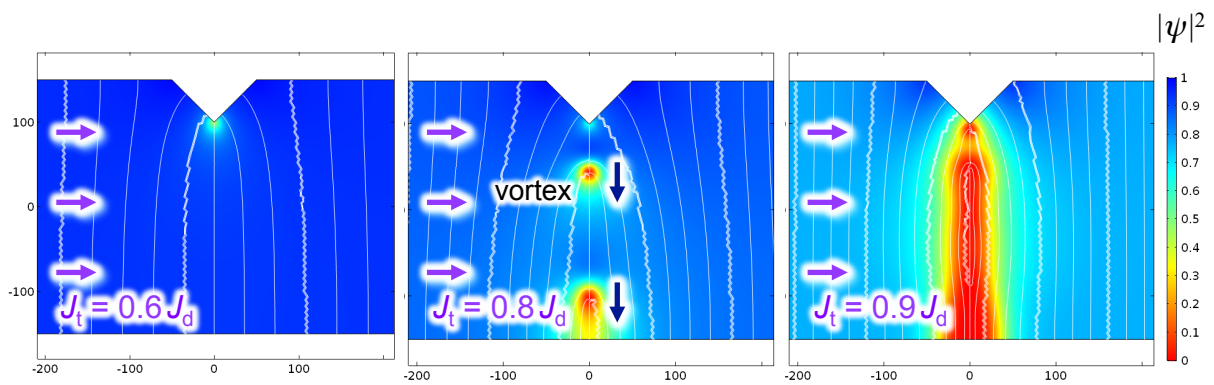
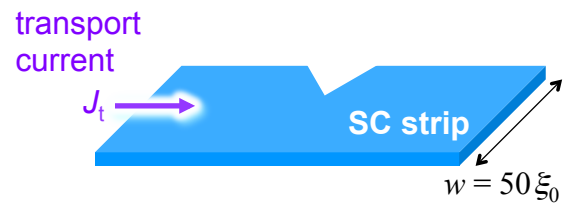
- vortex penetration



2D TDGL Simulation for Current Carrying SC Strips

- ◆ current-carrying SC strip with an edge defect (2D)

- vortex nucleation at the edge defect
- vortex flow



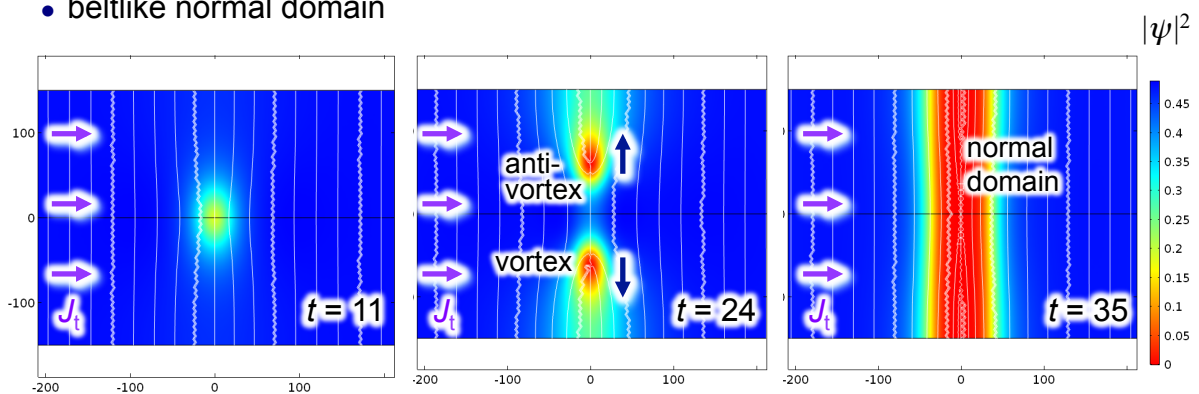
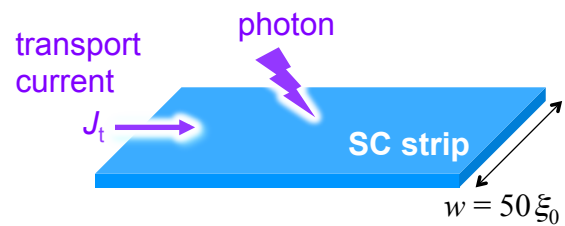
similar to the weak link (Josephson junction)

2D TDGL Simulation for SC Strip Photon Detectors

- ◆ SC strip photon detector (2D)

simulation based on the TDGL and heat diffusion equations

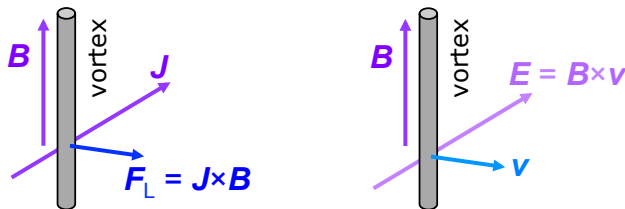
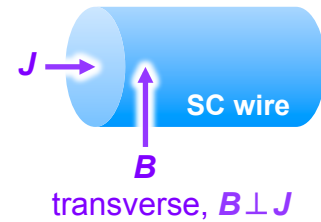
- hot spot due to the photon incidence
- vortex and antivortex nucleation
- vortex flow
- beltlike normal domain



Vortex Dynamics and Electromagnetic Response

◆ Vortex Dynamics in the Transverse ($B \perp J$) Configuration:

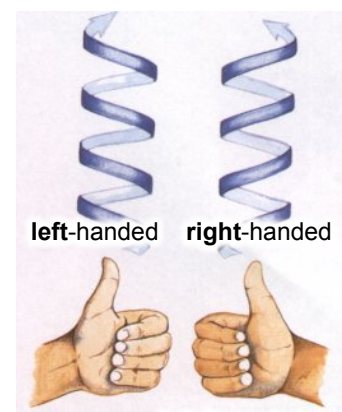
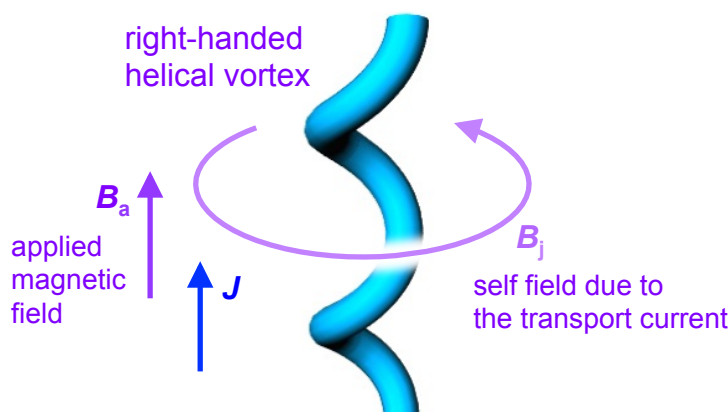
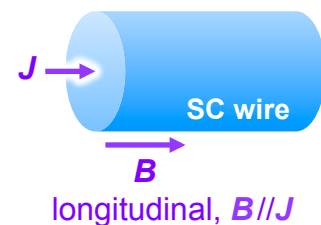
- driving (Lorentz) force on vortices, $F_L = J \times B$
 - ~ current density, J
- vortex velocity, v
 - ~ electric field, $E = B \times v$
- force balance equation for vortices, $F_L + F_p + F_v + \dots = 0$
 - (F_p : pinning force and F_v : viscous drag force)
 - ~ E - J characteristics, critical state model
 - ~ electromagnetic response of superconductors



Longitudinal Field Effect

◆ Vortex Dynamics in the Longitudinal ($B \parallel J$) Configuration:

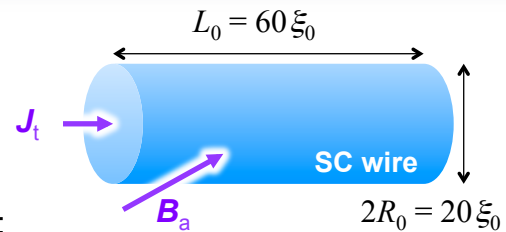
- $F_L = J \times B \sim 0$ (force free ?)
- v ? ~ $E \neq B \times v$?
- force balance equation ?
 - ~ E - J characteristics ??
 - (double critical state model,...)
 - ~ electromagnetic response ???
 - (e.g., large I_c , paramagnetic effect,...)



Lehninger et al.
"Principles of Biochemistry"

TDGL Simulation on Thin Superconducting Wires

- ◆ Superconducting cylinder:
 - radius $R_0 = 10 \xi_0$, length $L_0 = 60 \xi_0$
 - applied magnetic field $B_a = 0.2-0.4 B_{c2}$
 - transport current density $J_t = 0.3-0.7 J_d$



- ◆ TDGL equation for thin superconducting wires:
 - self field neglected, $A = B_a \times r / 2$

$$\begin{cases} \tau_0 \left(\frac{\partial}{\partial t} + i \frac{2\pi}{\phi_0} \Phi \right) \psi = \xi_0^2 \left(\nabla - i \frac{2\pi}{\phi_0} A \right)^2 \psi + (1 - |\psi|^2) \psi, \\ \nabla \cdot (\mathbf{j}_s + \mathbf{j}_n) = 0 \end{cases}$$

Parameters for $\text{YBa}_2\text{Cu}_3\text{O}_7$ at $T = 0\text{K}$

- τ_0 : GL relaxation time $\sim 0.03\text{ps}$
- ξ_0 : coherence length $\sim 1\text{nm}$
- λ_0 : penetration depth $\sim 100\text{nm}$
- B_{c2} : upper critical field $\sim 100\text{T}$
- J_d : depairing current density $\sim 10^{12} \text{A/m}^2$

- boundary conditions

$$\begin{cases} \mathbf{n} \cdot (\nabla - (2\pi i / \phi_0) A) \psi = 0 & \text{at the surface,} \\ \mathbf{n} \cdot \nabla \Phi = 0 & \text{at the side surface,} \\ \mathbf{z} \cdot \nabla \Phi = -J_t / \sigma_n & \text{at the ends (normal current injected)} \end{cases}$$

FEM Calculation by COMSOL Multiphysics®

Numerical TDGL-FEM simulation has been done by using the commercial software, **COMSOL Multiphysics®** (www.comsol.com).

- ◆ Coefficient Form PDE

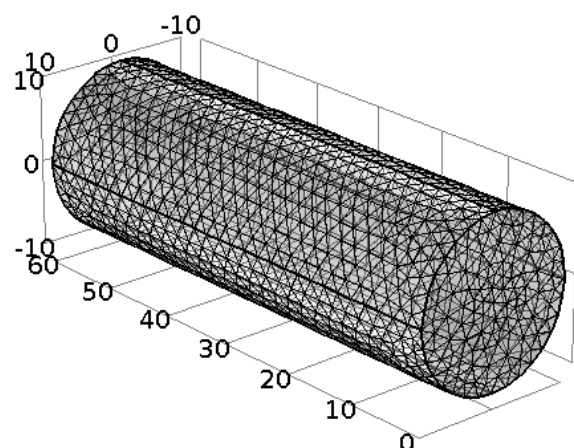
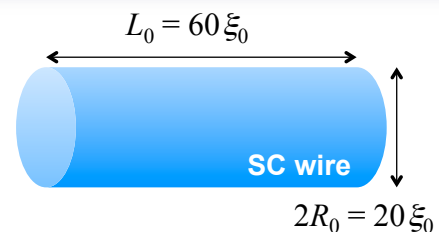
- order parameter $\psi(x, y, z, t)$
- scalar potential $\Phi(x, y, z, t)$

- ◆ Mesh:

- Element size parameter (in units of ξ_0):
 - max. element size 2.1,
 - min. element size 0.09
- Complete mesh consists of:
 - 33,317 domain elements,
 - 2,576 boundary elements, and
 - 180 edge elements

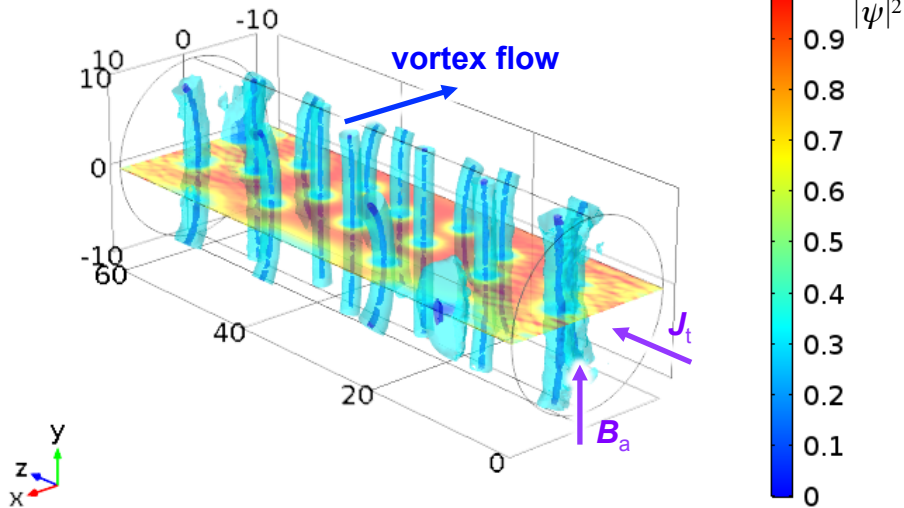
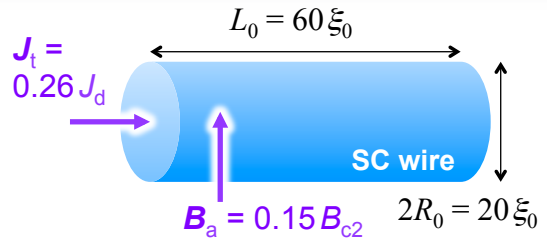
- ◆ Calculation:

- Number of degrees of freedom solved for 188,608 (plus 22,072 internal DOFs)
- Calculation time ~ 15 hours for $0 < t < 1,000$



Simulation Results (I): Transverse Field

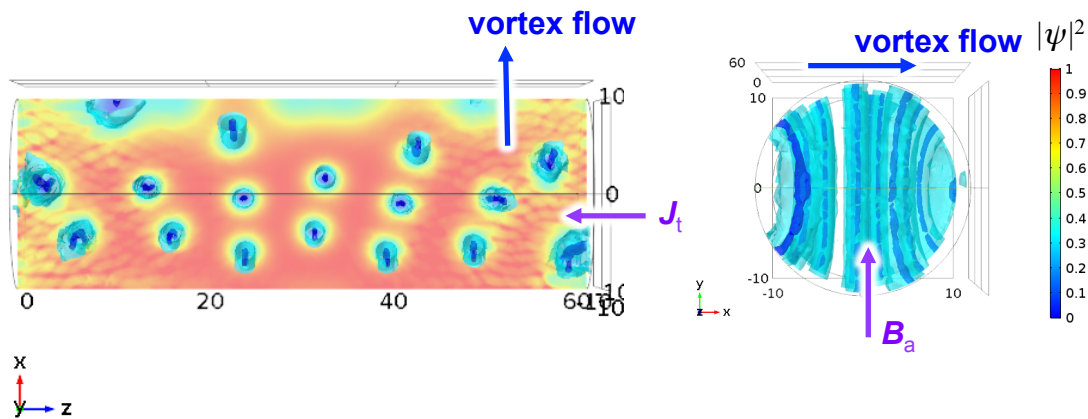
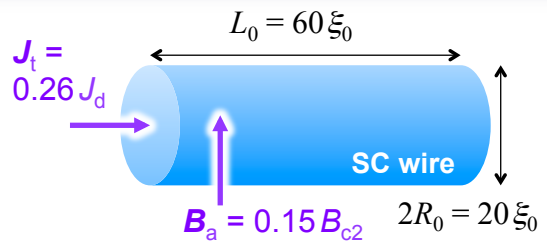
- ◆ Superconducting cylinder:
 - **transverse** magnetic field $B_a = 0.15 B_{c2}$
 - transport current density $J_t = 0.26 J_d$



Simulation Results (II): Transverse Field

- ◆ Superconducting cylinder:
 - **transverse** magnetic field $B_a = 0.15 B_{c2}$
 - transport current density $J_t = 0.26 J_d$

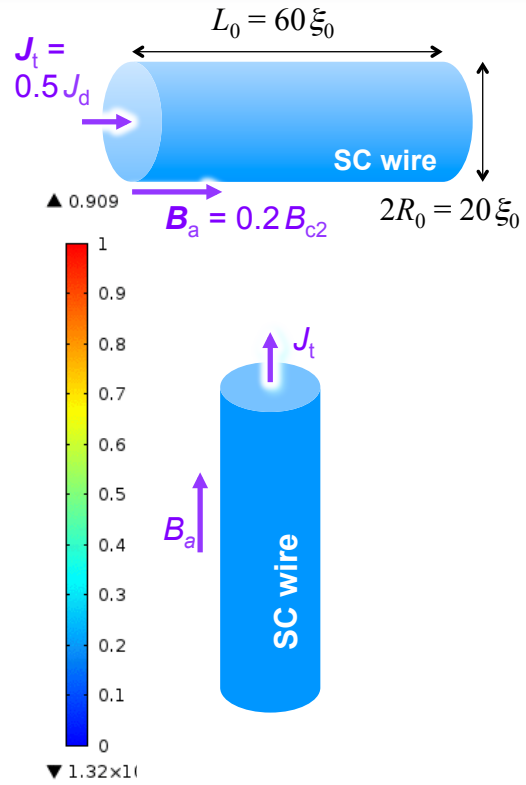
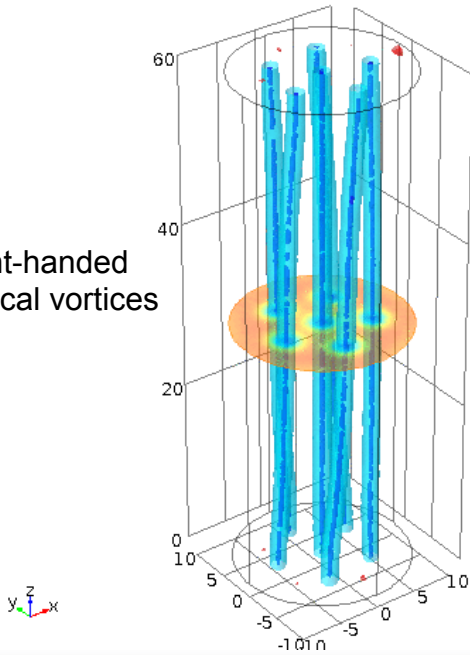
Time=170 s , $B_a/B_{c2} = 0.15$, $J_t/J_d = 0.26$



Simulation Results (III): Longitudinal Field

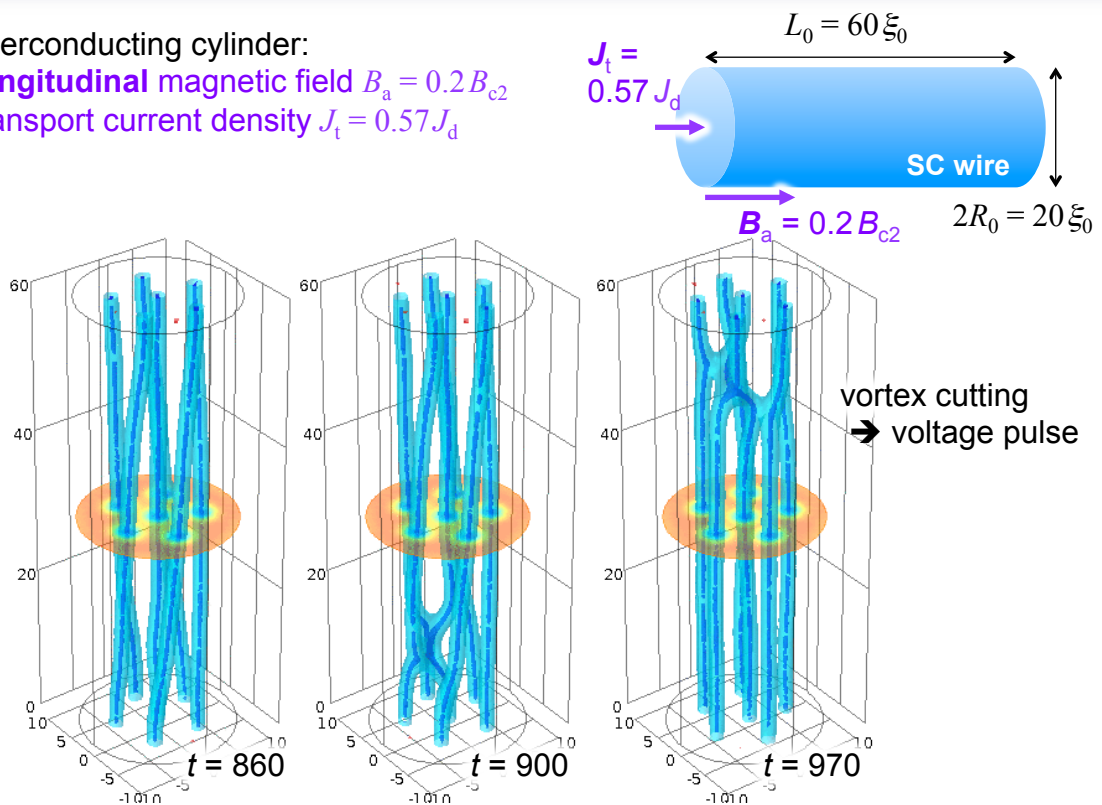
- ◆ Superconducting cylinder:
 - longitudinal magnetic field $B_a = 0.2 B_{c2}$
 - transport current density $J_t = 0.5 J_d$

- right-handed helical vortices



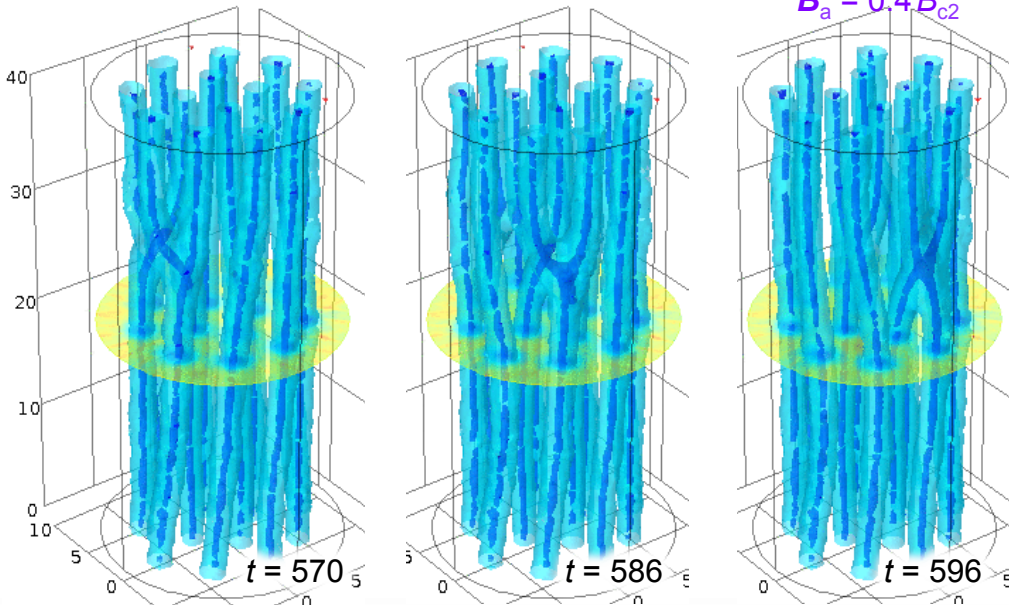
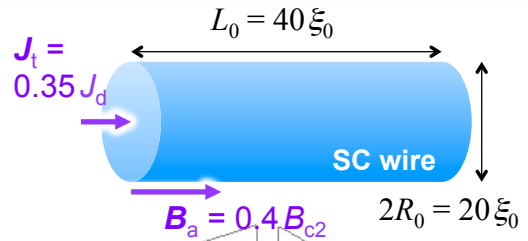
Simulation Results (IV): Longitudinal Field

- ◆ Superconducting cylinder:
 - longitudinal magnetic field $B_a = 0.2 B_{c2}$
 - transport current density $J_t = 0.57 J_d$



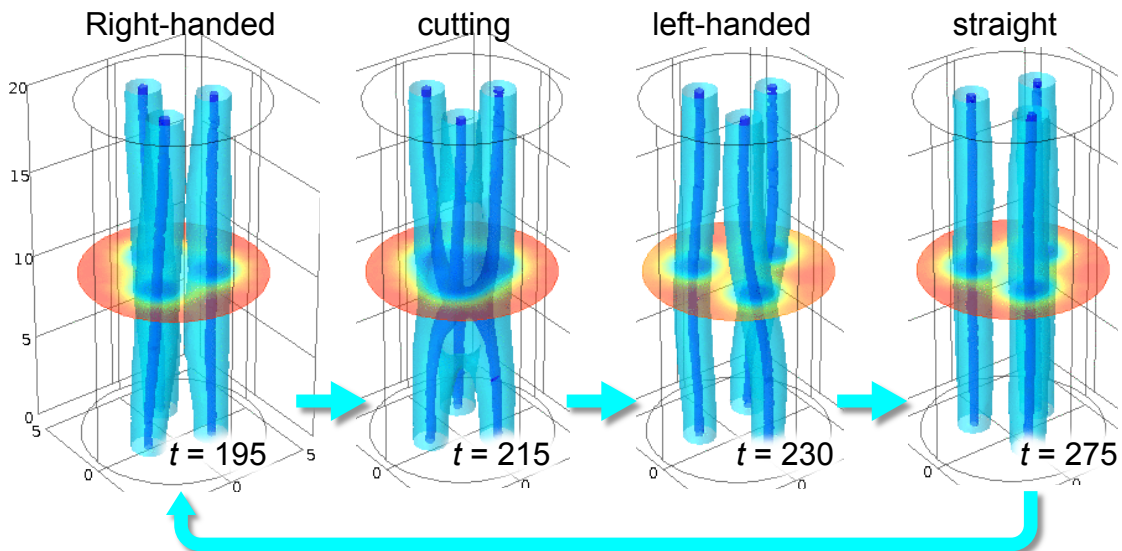
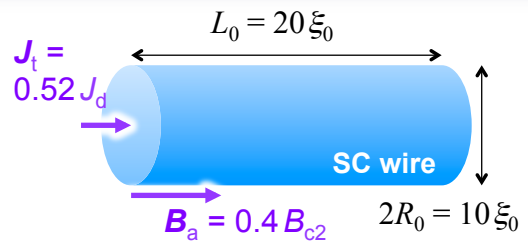
Simulation Results (V): Longitudinal Field

- ◆ Superconducting cylinder:
 - longitudinal magnetic field $B_a = 0.4B_{c2}$
 - transport current density $J_t = 0.35J_d$
 - propagation of vortex cutting



Simulation Results (VI): Longitudinal Field

- ◆ Superconducting cylinder:
 - longitudinal magnetic field $B_a = 0.4B_{c2}$
 - transport current density $J_t = 0.52J_d$



Summary

We have numerically investigated three-dimensional dynamics of vortices in thin superconducting (SC) cylinders exposed to transport currents and longitudinal magnetic fields.

- The time-dependent Ginzburg-Landau (TDGL) equations are useful for modelling of nanoscale (nm– μm) SC phenomena (e.g., vortex dynamics, SC electronics devices,...).
- Helical vortex structure in current carrying SC wires in longitudinal magnetic fields
- Complicated vortex dynamics with cutting for large current:
 - propagation of vortex cutting in many vortices
 - right-handed helices \rightarrow cutting \rightarrow left-handed helices \rightarrow ...

Future issues:

- Further investigation (by changing applied field, current, and sample size,...) to clarify the complicated vortex dynamics
- Large scale simulation for bulk behavior
- Pinning, anisotropy, thermal fluctuation,...