

## Eddy current modeling in multifilamentary superconductive tapes

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# Outline

- Context
- Modeling approach & Formulation
- Numerical examples

# Context

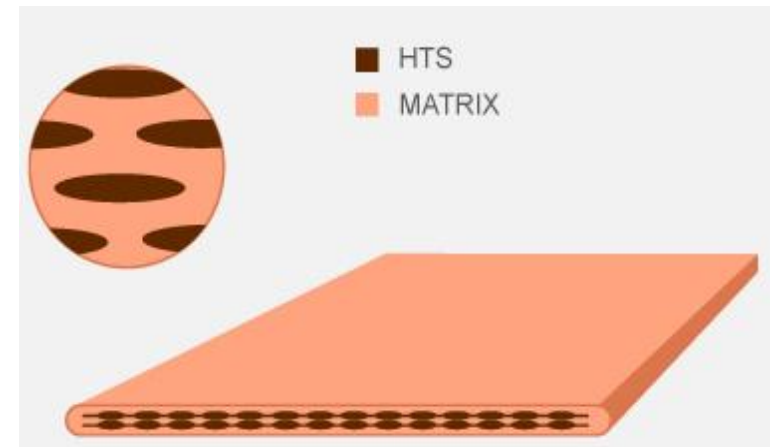
Modeling of the ELM filed interaction with composites superconductive tapes

## Challenges

- Complex geometries (multiscale dimensions)
- Material proprieties (Nonlinearity , anisotropy)
- Limits of conventional numerical tools

## Motivations

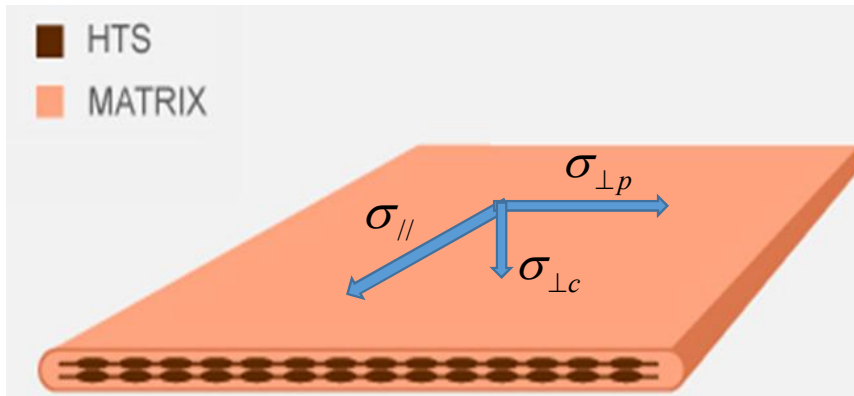
Rapid eddy current losses modeling in multifilamentary HTS composites tapes



<http://www.superox.ru/en/superconductivity/>

# Approach

Multi-filamentary composite tapes



$$\bar{\sigma}(J) = \begin{bmatrix} \sigma_{//}(J) & 0 & 0 \\ 0 & \sigma_{\perp p} & 0 \\ 0 & 0 & \sigma_{\perp c} \end{bmatrix}$$

$$\mu = \mu_0$$

$$\sigma_{//} = \eta_s \frac{J_{//}}{E_{//}} + (1 - \eta_s) \sigma_m$$

$$E_{//} = E_c \left( \frac{J_{//}}{J_c(B)} \right)^n$$

$\sigma_{//}$ ,  $\sigma_{\perp p,c}$  depend on:

- The HTS filaments volume fraction ( $\eta_s$ ) and shape
- The matrix conductivity
- The interface between the filaments and matrix
- The source field orientation

$$J_c(B) = \frac{J_{c0}}{1 + \frac{|B|}{B_0}}$$

# Formulation

$$\vec{\nabla} \cdot \vec{J} = 0 \rightarrow \vec{J} = \vec{\nabla} \times \vec{T} \quad \vec{\nabla} \times \overline{\overline{\rho}} \vec{\nabla} \times \vec{T} = -\partial_t \vec{B} \quad (1)$$

$$\vec{B} = \vec{B}^s + \vec{B}^r = \underbrace{\vec{B}^{ext}}_{\vec{B}^s} + \mu_0 \overline{\overline{M}} \vec{J}_s + \mu_0 \overline{\overline{M}} \overline{\overline{C}} \vec{T} \quad (2)$$

$$\overline{\overline{C}} \overline{\overline{R}}(J) \overline{\overline{C}} \vec{T} = -i\omega \left( \vec{B}^s + \mu_0 \overline{\overline{M}} \overline{\overline{C}} \vec{T} \right) \quad (3)$$

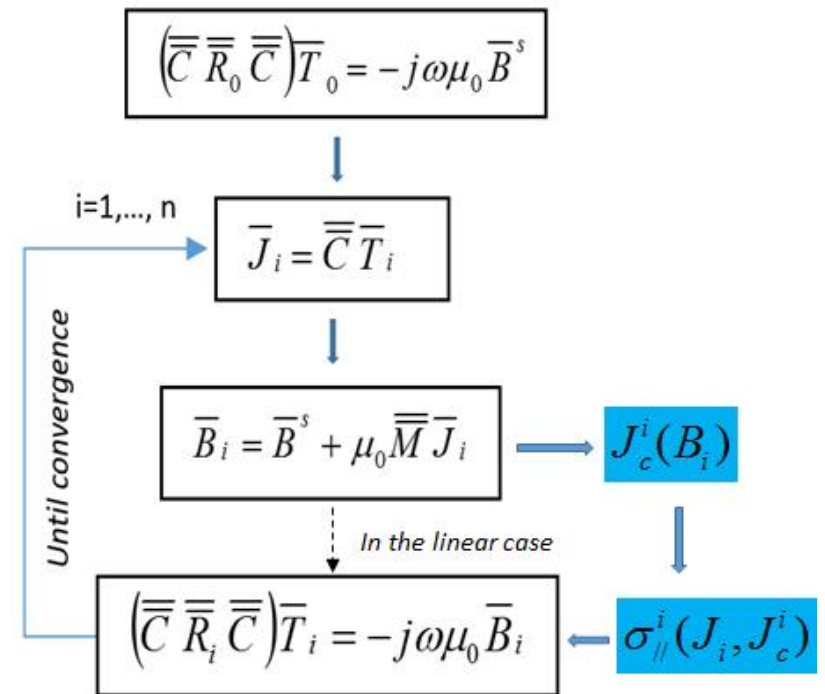
Thin layer approximation  $\mathbf{J}=(J_x, J_y, 0)$ ,  $\mathbf{B}=(0, 0, B_z)$

$$\overline{\overline{C}}_{x,y} \overline{\overline{R}}_{x,y}(J) \overline{\overline{C}}_{x,y} \vec{T}_z = -i\omega \left( \vec{B}_z + \mu_0 \overline{\overline{M}}_{x,y} \overline{\overline{C}}_{x,y} \vec{T}_z \right) \quad (3-bis)$$

Only the **active parts** are discretized

The problem is solved **iteratively** in the frequency domain

$\overline{\overline{C}}$  Curl matrix  
 $\overline{\overline{M}}$  Integral matrix  
 $\overline{\overline{R}}$  Resistivity matrix



# Numerical example

Typical Bi-2223 tape

$$W_t = 3 \cdot 10^{-3} \text{ m}$$

$$T_t = 2 \cdot 10^{-4} \text{ m}$$

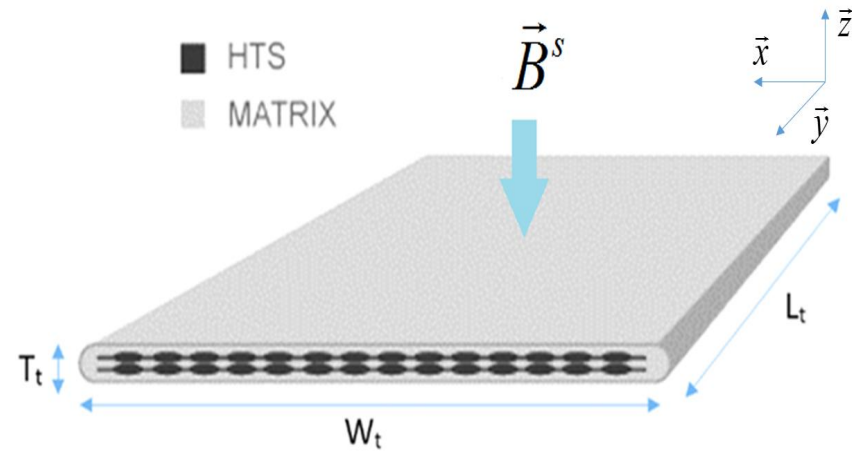
$$L_t > L_C$$

$$\eta_s = 0.5$$

$$J_{c0} = 2 \cdot 10^7 \text{ A/m}^2$$

$$n = 25$$

$$\rho_m = 16 \cdot 10^{-9} \text{ }\Omega\cdot\text{m}$$

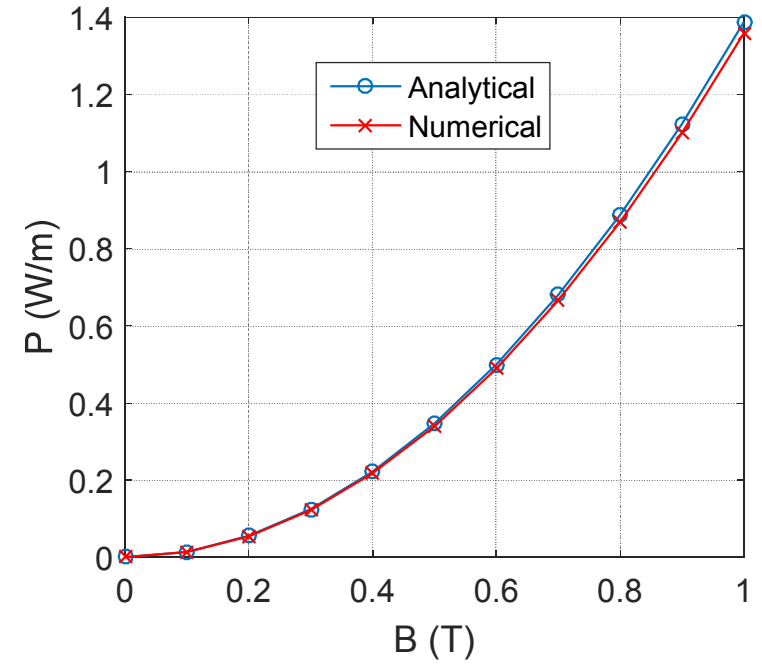
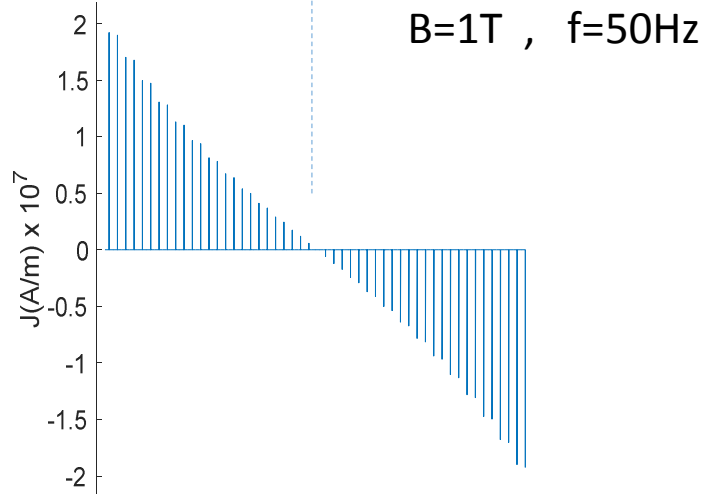
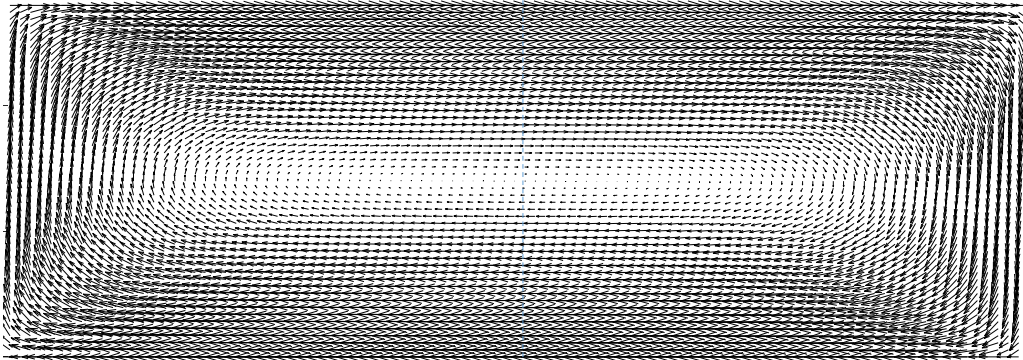


$$J_c(B) = \frac{J_{c0}}{1 + 3.07|B|}$$

From M. P. Oomen (2000)

## Validation

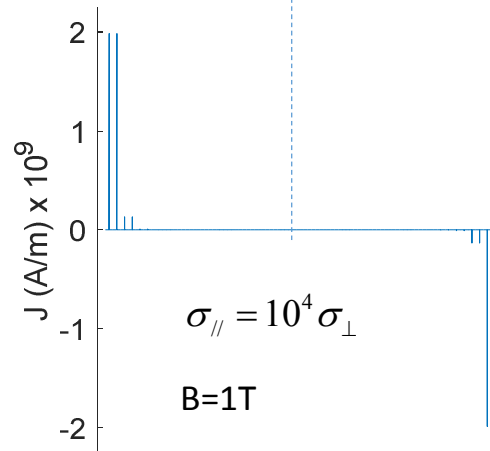
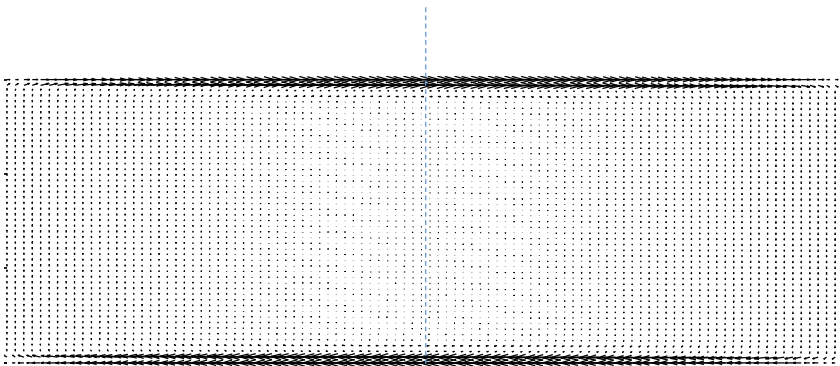
$$\sigma_{//} = \sigma_{\perp}$$



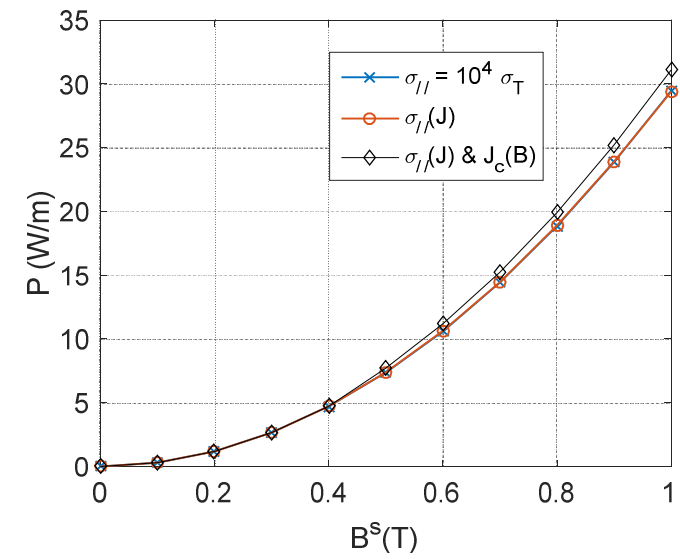
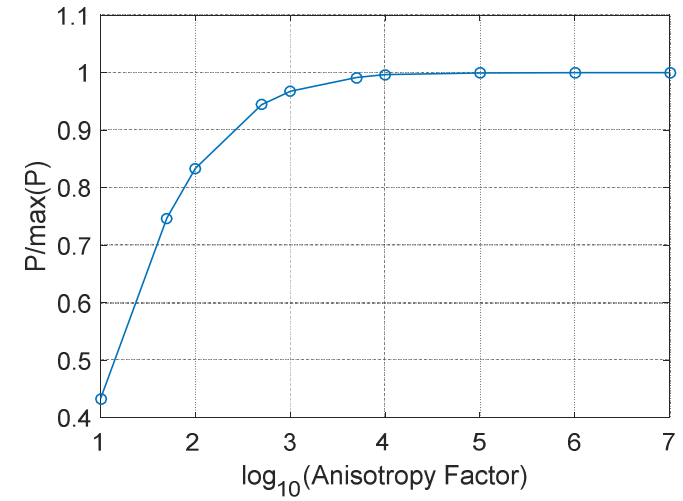
Static error of 2%

\*Analytical model from M. P. Oomen (2000)

$$\sigma_{//} \gg \sigma_{\perp}$$

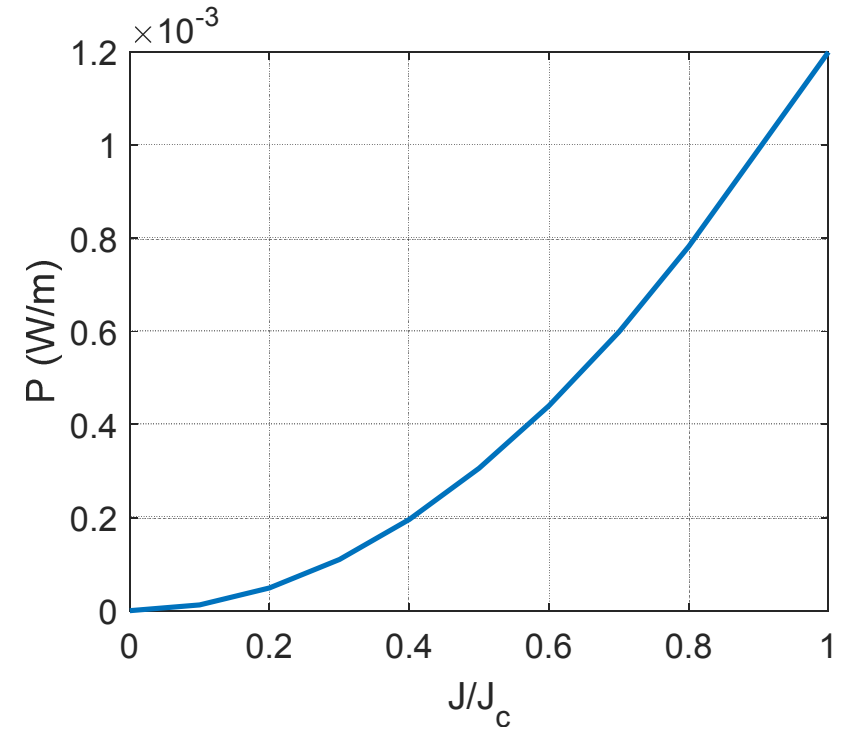
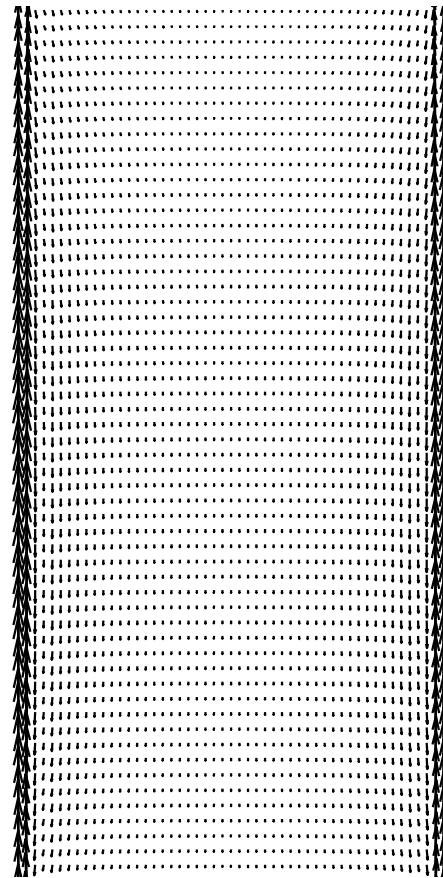
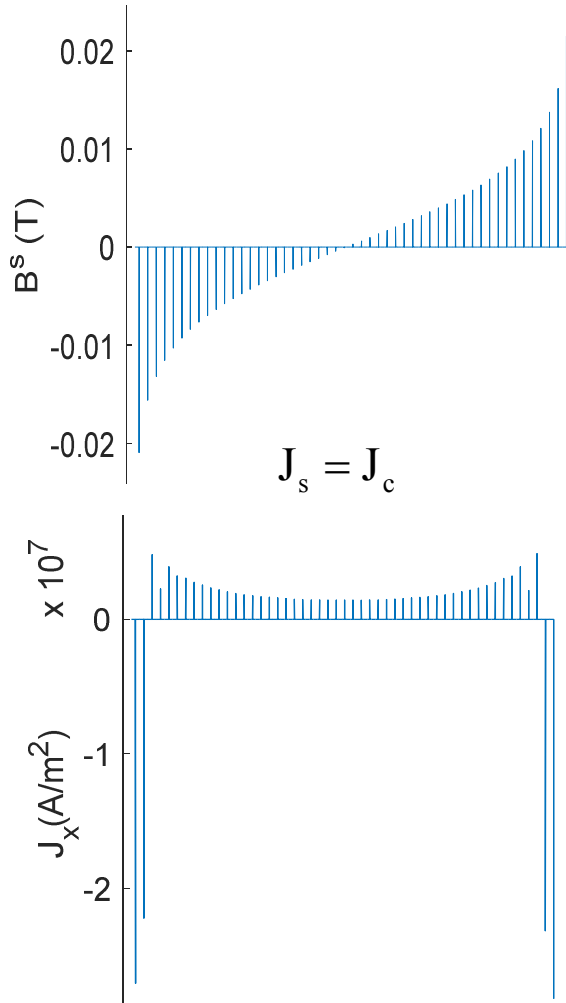


f=50Hz





### AC losses due to transport current



$$\sigma_{//} = 10^4 \sigma_{\perp}$$

$$f = 50 \text{ Hz}$$

$$J_c = 1.16 \cdot 10^8 \text{ A/m}^2$$

## Conclusions

- Rapid approaches are needed for superconductive material modeling
- Better compromise between the accuracy and computing time
- Meet the progress in numerical tools (storage capacity)
- Integral approaches are promising
- Simplifications can be made on material behavior and structures in some cases

**Thank you for your attention**